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Journal Article peer-reviewed accepted version (Postprint) as:

DOI of this document\*(secondary publication): https://doi.org/10.26092/elib/3677 Publication date of this document:

17/02/2025

\* for better findability or for reliable citation

#### Recommended Citation (primary publication/Version of Record) incl. DOI:

Twist and Return-Induced Ring Strain Triggers Quick Relaxation of a (Z)-Stabilized Cyclobisazobenzene Chavdar Slavov, Chong Yang, Andreas H. Heindl, Tim Stauch, Hermann A. Wegner, Andreas Dreuw, and Josef Wachtveitl. The Journal of Physical Chemistry Letters 2018 9 (16), 4776-4781 DOI: 10.1021/acs.jpclett.8b02159

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# Twist and Return – Induced Ring Strain Triggers Quick Relaxation of a (*Z*)-Stabilized Cyclobisazobenzene

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**ABSRTRACT**: Continuous irradiation of the thermodynamically stable (Z,Z)cyclobisazobenzene does not lead to accumulation of a (Z,E) or (E,E) isomer as one might expect. Our combined experimental and computational investigation reveals that  $Z \rightarrow E$ photoisomerization still takes place on an ultrafast timescale, but induces large ring strain in the macrocycle, which leads to a very fast thermal back-isomerization preventing photo-stationary accumulation of (E)-isomers.

**TOC GRAPHICS** 



Azobenzene (AB) derivatives represent an important family of molecular photoswitches that are extensively used for spatiotemporally precise and reversible control of a variety of nanostructures<sup>1-3</sup> and reactions<sup>4-5</sup> of interest to chemistry and biology. Furthermore, ABs are used as building blocks in larger aggregates to form photoresponsive materials.<sup>6-9</sup> One of the main application advantages of AB photoswitches is the substantial geometrical change (molecular length change of ~3.4 Å) associated with the  $E \leftrightarrow Z$  photoinduced isomerization of the central N=N bond. In addition, various AB derivatives can be prepared efficiently using a number of systematic methods.<sup>10</sup> ABs also show a remarkable fatigue resistance.<sup>11</sup>

The progress in the development of advanced photoresponsive nanostructures and materials, where multiple AB units are linked and required to operate cooperatively, demonstrates that the current understanding of the photochemistry of single ABs does not suffice to predict the behavior of more complex AB constructs. In such multi-AB constructs, the interaction between the individual photoswitching units leads to versatile effects like excitonic coupling,<sup>12-14</sup> extended  $\pi$ -delocalization,<sup>14-16</sup>  $\pi$ -stacking,<sup>9,17</sup> steric hindrance<sup>18-19</sup>, and cooperativity<sup>8,20</sup>, which all can strongly modulate the physicochemical properties of the AB units. Therefore, the rational design of photoresponsive systems requires prior, detailed investigation of all these diverse phenomena and a thorough understanding of their complexity.

Photoswitchable macrocyclic ABs<sup>21</sup> are a family of complex, cyclic oligo-AB constructs with great potential for application in fields like host-guest interaction,<sup>22-23</sup> supramolecular chemistry,<sup>9,17,24</sup> molecular machines,<sup>13</sup> and energy storage<sup>25</sup>. In oligo-AB macrocycles, steric distortions and ring strain alter the photoisomerization dynamics and the thermal relaxation rate of the individual AB units, and thus impact the overall photoisomerization quantum yield (QY).<sup>22,26-27</sup> Nevertheless, the remarkable photoswitching property of the individual ABs is

preserved in this type of compounds. However, a unique case of a minimal, strongly constraint cyclotrisazobenzene (CTA) that is locked in its (*E*,*E*,*E*)-isomer has been reported recently.<sup>28-29</sup> Indeed, the geometric constraints in CTA (only minimal changes of  $\ll$ CNNC are possible) transform the well-known, efficient S<sub>1</sub>–S<sub>0</sub> conical intersection (CI), which typically drives the photoisomerization of ABs, into an ultrafast (<1 ps), non-photochemical relaxation pathway back to the initial isomer.<sup>29</sup>

To date, only a few reports exist on AB derivatives for which the (*Z*)-isomer is the thermodynamically stable form. For example in 5,6-dihydrodibenzo[c,g][1,2]diazocine,<sup>30-31</sup> in [0.0](3,3')-azobenzenophane,<sup>32</sup> and in [1.1](4,4')-azobenzenophane<sup>33</sup>. The first compound is a monomeric, bridged AB, while in the other two compounds, the (*Z*)-stable AB units are part of sterically hindered cyclic bisABs. Interestingly,  $Z \rightarrow E$  photoisomerization was reported only for the former two compounds, but no (*E*)-isomer could be accumulated for [1.1](4,4')-azobenzenophane (see chemical structure in Fig. 1), which raises the question whether [1.1](4,4')-azobenzenophane is capable of isomerization at all or it is locked like the previously reported CTA<sup>29</sup>. Moreover, in the case of the two cyclic bisABs, an in-depth understanding of the molecular origin of the observed properties is still lacking. This motivated us to investigate the photoinduced dynamics of [1.1](4,4')-azobenzenophane cyclobisazobenzene (CBA) by combining ultrafast optical spectroscopy and theoretical analysis. More fundamentally, our work contributes to the understanding of the AB photoisomerization dynamics as it reports on the ultrafast dynamics of a unique thermodynamically stable (*Z*)-AB.

The stationary absorption spectrum of the thermodynamically stable (*Z*,*Z*)-isomer of CBA possesses the typical (*Z*)-AB  $\pi\pi^*$ -transition (292 nm) and n $\pi^*$ -transition (430 nm) bands (Fig. 1). Interestingly, the intensity of the n $\pi^*$ -transition band appears to be slightly amplified in CBA as

compared to a spectral superposition of the  $n\pi^*$  bands of two independent monomeric (*Z*)-AB molecules. Continuous irradiation at the absorption maxima of CBA with high power LEDs or with a HgXe arc lamp did not induce any spectral changes, and thus a stationary detection of CBA  $Z \rightarrow E$  photoswitching was not feasible.



**Figure 1.** Experimental absorption spectrum of cyclobisazobenzene (structure shown as an inset) in ethanol. The vertical bars indicate the excitation wavelengths (304 nm and 437 nm) used in the corresponding time-resolved experiments.

Ultrafast transient absorption (TA) experiments (see Ref. [<sup>34</sup>] and the SI for a description of the experimental setup) were performed to investigate the photoinduced dynamics of (*Z*,*Z*)-CBA and to clarify whether the AB units of CBA have retained their isomerization capability. The TA data after excitation in the  $\pi\pi^*$ -transition band (304 nm) and after excitation in the  $\pi\pi^*$ -transition band (304 nm) and after excitation in the  $\pi\pi^*$ -transition band (437 nm) show identical dynamics (Fig. 2A & 2C, see also Fig. S1), and both, indeed, reveal significant formation of the (*E*)-isomer as photoproduct. This is concluded from: *i*) the non-decaying  $\pi\pi^*$ -band ground state bleach signal (GSB) of the (*Z*)-isomer (410-450 nm, present

>1 ns); and *ii*) the absorption signals associated with the  $\pi\pi^*$  and the  $n\pi^*$  bands of the formed (*E*)-isomer (below 400 nm and above 450 nm). The major difference between the two datasets is found on the sub-600 fs timescale. While after  $\pi\pi^*$  excitation the sub-600 fs TA data (Fig. 2A) is dominated by excited state absorption (ESA) over the complete spectral range, after  $n\pi^*$  excitation the data in the 425-475 nm spectral range show a nearly instantaneous GSB of the  $n\pi^*$  band. Consequently, the sub-600 fs ESA contribution and its dynamics in the  $\pi\pi^*$  excitation experiment can be straightforwardly assigned to the (*Z*,*Z*)-CBA's  $\pi\pi^*$  excited state (conventionally labeled as S<sub>2</sub> in monomeric ABs). The TA data also show significant spectral overlap between the ESA bands and the (*E*)-isomer product bands, which largely conceals the early excited state relaxation, since the decaying ESA signatures transform into product absorption signatures. Nevertheless, the product absorption bands, located in the spectral ranges below 400 nm and above 450 nm, are clearly seen already <1 ps indicating ultrafast (*E*)-isomer formation after both  $n\pi^*$  and  $\pi\pi^*$  excitation of (*Z*,*Z*)-CBA.

Further insight into the early photoisomerization dynamics of (Z,Z)-CBA is gained via the corresponding lifetime density maps (LDM) (Fig. 2B & D), obtained from the lifetime distribution analysis of the TA data (see Ref. [<sup>35</sup>] and SI for technical details). After  $n\pi^*$  excitation (437 nm), the LDM (Fig. 2D) shows a series of lifetime distributions centered at ~200 fs and associated with: i) the decay of the ESA, visible at the edges of the detection spectral range; ii) the partial recovery of the GSB (400-470 nm range); and iii) the formation of the vibrationally hot (*E*)-AB product absorption bands (below 400 nm and above 450 nm). Therefore, the major decay of the  $n\pi^*$  excited state of (*Z*,*Z*)-CBA and the onset of the  $Z \rightarrow E$  photoisomerization reaction occur with ~200 fs lifetime, which is similar to what was observed for monomeric (*Z*)-AB<sup>14,30,36-37</sup>. In previous studies on (*Z*)-AB, an additional 1-3 ps lifetime

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component has been reported and assigned to diffusive motion on the  $n\pi^*$  PES<sup>14,36-37</sup> or alternatively to the decay of a dark excited state intermediate<sup>38</sup>. However, such a component appears to have only a minor contribution to the LDM of (*Z*,*Z*)-CBA (Fig. 2D), where at ~1 ps mainly the onset of the non-exponential ground state vibrational cooling dynamics<sup>36,39</sup> is detected as elongated and tilted lifetime distributions extending towards 10-30 ps.



**Figure 2.** A) and C) Transient absorption data of the (Z,Z)-CBA measured in ethanol after 304 nm and 437 nm excitation, correspondingly (ground state bleach – negative absorption difference signal, light to dark blue; excited state and hot ground state absorption – positive absorption difference signal, yellow through red to black); B) and D) Corresponding lifetime density maps (LDMs) obtained from the lifetime distribution analysis (see SI) of the time-resolved data in A) and C). The reading of the LDMs is as for a DAS from global lifetime analysis – i) positive (red) amplitudes account for decay of absorption or rise of GSB; ii) negative (blue) amplitudes account for rise of absorption or decay of GSB.

In contrast, the LDM from the 304 nm excitation experiment (Fig. 2B) shows a pronounced positive amplitude contribution with a lifetime of ~200 fs, accounting for the decay of the  $\pi\pi^*$ excited state, comparable to what has been reported for monomeric (Z)-AB.<sup>38,40</sup> The similar decay lifetimes of the  $\pi\pi^*$  and the  $n\pi^*$  excited states in (Z,Z)-CBA does not permit their clear distinction in the  $\pi\pi^*$ -excitation experiment. Nevertheless, after the main ~200 fs decay, two minor negative lifetime distributions can be seen at  $\sim$ 1-2 ps, which, as described above, account for the onset of the ground state cooling commencing with the end of the relaxation of the  $n\pi^*$ state to the ground state. The lifetime distributions, which account well for such non-exponential behavior, reflect this situation in a more natural way. The cooling of the vibrationally hot CBA molecules in the ground state is accomplished with the characteristic 10-30 ps lifetime<sup>36,39</sup>, as reflected by the elongated and spectrally tilted lifetime distributions on this scale (Fig. 2B & D). Additional,  $\sim$ 500 ps lifetime distributions (negative amplitude) are detected in the (E)-isomer absorption ranges (below 400 nm and above 450 nm) in both LDMs. These distributions can be linked to some conformational dynamics in the ground state that causes some modulation of the product absorption bands.

Although (*E*)-isomer of CBA could not be accumulated under stationary irradiation, the femtosecond experiments show clearly that CBA undergoes  $E \rightarrow Z$  isomerization with significant quantum efficiency. Therefore, flash photolysis experiments were performed to determine the thermal  $E \rightarrow Z$  relaxation of CBA, which yielded a 7.7 ms half-life for the (*E*)-isomer (Fig. S2). This high relaxation rate indicates that the formed (*E*)-configuration is significantly distorted and unstable in the electronic ground state, and therefore explains why the molecule cannot retain it.

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To elucidate the molecular mechanisms of the observed photoisomerization dynamics of CBA, quantum chemical calculations were performed using time-dependent density functional theory<sup>41-42</sup> in combination with the BHLYP functional,<sup>43</sup> 6-31G\* basis set.<sup>44</sup> and a conductorlike polarizable continuum model for ethanol solvation ( $\varepsilon = 24.5$ )<sup>45</sup> (see SI for details). Ground state geometry optimization of CBA revealed two possible geometries (see, Fig. S3 and Table S1) for each of the three isomerization states [(Z,Z), (Z,E), (E,E)] of CBA. These two geometries, however, are very similar and have negligible energy difference. The structural parameters (bond lengths and dihedral angles) of (Z,Z)-CBA closely resemble those of (Z)-AB (Tables S1), indicating that the thermodynamically stable isomer of CBA is only marginally, distorted compared to the isolated (Z)-AB. In contrast, the dihedral angles of the (E)-subunits in (Z,E)- or (E,E)-CBA are significantly affected and differ from undistorted (E)-AB. For example, the largest distortion of  $\ll$ CNNC is detected for the (E)-subunit in (Z.E)-CBA (146.6°), while in (E,E)-CBA both (E)-subunits are closer to being planar (157.3°) (Table S1). This explains why the (Z,E)-isomer of CBA has higher total energy than the (E,E)-isomer. A similar tendency is detected in the strain analysis<sup>46-48</sup> of CBA (see SI for details), which shows that there is hardly any strain energy (Table S2, Fig. S7) stored in the (Z)-subunits of both (Z,Z)-CBA and (Z,E)-CBA (the energy is similar to that of monomeric (Z)-AB). Conversely, tremendous molecular strain is stored in the distorted (*E*)-subunits of both (*Z*,*E*)-CBA and (*E*,*E*)-CBA – particularly in the  $\ll$ CNNC (21-30.1%), while the rest is delocalized over the benzene rings (Fig. S7). In the (E,E)-CBA the strain energy is shared between the two (E)-subunits (each stores  $\Delta E_{harm}$  =  $\sim$ 17.7 kcal/mol, as calculated with the JEDI analysis on the basis of harmonic approximation<sup>49-</sup> <sup>50</sup>), however, in the (Z,E)-CBA this energy is almost exclusively stored in the single (E)-subunit, which is therefore particularly strained ( $\Delta E_{harm} = ~27 \text{ kcal/mol}$ ) (see Table S2). The large amount

of strain in the (*E*)-subunits is the driving force of the high thermal relaxation rate (half-life 7.7 ms) observed in the flash photolysis experiments (Fig. S2) and the inability to accumulate the (*E*)-isomers under continuous irradiation.

The simulated spectra (Fig. S4) and the attachment/detachment densities (Fig. S5) show that the absorption band at 430 nm (Fig. 1) should be a superposition of the nearly degenerate  $n\pi^*$ transitions ( $S_1$  and  $S_2$ ) of the two (Z)-subunits of (Z,Z)-CBA. The observed slightly increased amplitude of the 430 nm absorption band (Fig. 1) indicates that the structural distortion in CBA may lead to an amplified vibronic intensity borrowing from the  $\pi\pi^*$  to the  $n\pi^*$  transitions. This interpretation is also supported by the optimized ground state structures of the (Z,Z)-isomers (Fig. S3) and the crystal structure of the molecule<sup>33</sup>, which show that the benzene rings in CBA are slightly tilted as compared to the structure of monomeric (Z)-AB. In effect, the angle between the planes of the lone pair of the N-atoms and the p-orbitals of the benzene rings (see the ∢CNNC dihedral differences in Table S1) is reduced, thus possibly facilitating intensity borrowing. The higher energy transitions (S<sub>3</sub> and S<sub>4</sub>) of CBA are of  $\pi\pi^*$  character (Fig. S5) and in the case of (Z,Z)-CBA account for the 304 nm band in the experimental absorption spectrum (Fig 1 and Fig. S4).  $Z \rightarrow E$  isomerization of one AB unit leads to a red shift of the associated  $n\pi^*$ transition and an increased intensity (Fig. S4) due to the non-planar geometry of the (E)-subunit and the ensuing breaking of the  $C_{2h}$  symmetry. We observe a stronger red shift of the  $n\pi^*$ transition band for one of the possible (Z,E)-CBA ground state geometries apparently due to different interactions of the lone-pairs with the remaining nuclear frame (Fig. S3 & Fig. S4). The  $Z \rightarrow E$  isomerization also induces a red shift of the  $\pi\pi^*$  transition of the (E)-isomer (Fig. S4) in agreement with the experimental data (Fig. 2 and Fig. S2).

To understand the unique isomerization behavior of CBA, relaxed potential energy surface (PES) scans along the rotational coordinate were performed in the S<sub>1</sub> ( $n\pi^*$ ) and S<sub>3</sub> ( $\pi\pi^*$ ) states of CBA (Fig. 3). The shape of the  $S_1$  PES of both azobenzene units in CBA resembles that of the monomeric AB (cf. [51-52]). We do not observe an energetic barrier on the relaxation pathway towards the conical intersection (CI) with the ground state (Fig. 3). Furthermore, the close energetic proximity of the S<sub>1</sub> and the S<sub>0</sub> state at  $\ll$ CNNC of 80° facilitates the ultrafast excited state decay through the CI leading to isomerization via the S<sub>1</sub>-S<sub>0</sub> seam already in the vicinity of this CI. On the ground state PES of the (Z,E)-isomer (middle of Fig. 3A), the molecular conformations towards isomerization of the remaining (Z)-subunit are energetically uphill. Therefore, the (Z, E)-isomers would predominantly reside in the local minimum on the left side of the vertical dashed line in Fig. 3A. Excitation of these populations would lead to backisomerization since the  $S_1$  PES shows a steeper slope towards the (Z,Z)-isomer, favoring the  $E \rightarrow Z$  isomerization direction. If the (Z,E)-population in the ground state is thermally shifted towards the conformations with increasing  $\ll$ CNNC for the remaining (Z)-subunit, then excitation to the left side of the maximum in the S<sub>1</sub> PES could potentially lead to a second  $Z \rightarrow E$ isomerization and formation of the (E,E)-isomer. However, this has not been observed so far. Concerted isomerization of the two (Z)-subunits of CBA is highly unlikely due to high energetic barriers on all excited states PESs in the direction of concerted  $ZZ \rightarrow EE$  isomerization (see Fig. S10).



**Figure 3.** Potential energy surfaces for the S<sub>1</sub> (A) and S<sub>3</sub> (B) optimized geometries with constrained  $\ll$ CNNC dihedral angles at the TDDFT/BHLYP/6-31G\* level of theory (For S<sub>0</sub> optimized PES see Fig. S8). The figures describe the photoisomerization dynamics after n $\pi^*$  (A) and  $\pi\pi^*$  (B) excitation. In (A) the vertical dashed line indicates the point where the constriction of  $\ll$ CNNC is switched from the first to the second AB-unit, therefore, the possibility of consecutive isomerization from the (*Z*,*Z*)-isomer to the (*E*,*E*)-isomer after n $\pi^*$  excitation is also explored. CI – conical intersection; IC – internal conversion.

The relaxed scans along the  $\ll$ CNNC rotation coordinate for the PES of the S<sub>3</sub> ( $\pi\pi^*$ ) optimized geometry of CBA show an uphill slope toward 90° and the presence of an S<sub>3</sub>-S<sub>2</sub> crossing point at ~60° of  $\ll$ CNNC (Fig. 3B). Furthermore, the S<sub>3</sub> and S<sub>2</sub> PESs lie energetically close together (see also Fig. S9). Therefore, excitation into the  $\pi\pi^*$  band (maximum at ~292 nm) would lead to a rapid S<sub>3</sub> relaxation to the S<sub>2</sub> (n $\pi^*$ ) PES and further to the S<sub>1</sub> (n $\pi^*$ ) PES. The subsequent S<sub>1</sub> state relaxation then proceeds via the pathway described above for n $\pi^*$  excitation (*cf.* Fig. 3A).

In conclusion, our combined experimental and theoretical study of CBA, a model macrocyclic AB, demonstrates that the individual AB units still undergo ultrafast  $Z \rightarrow E$  photoisomerization despite the strong steric hindrance in the molecule. However, the internal molecular strain, induced by the  $Z \rightarrow E$  isomerization, leads to a high relaxation rate, and therefore prevents recognizable accumulation of the (Z,E)-isomer under continuous irradiation. Interestingly, the steric distortion in CBA, as compared to monomeric AB, breaks the molecular symmetry, which gives rise to an increased intensity of the  $n\pi^*$  transition due to vibronic intensity borrowing. The ultrafast dynamics, as revealed by the transient absorption experiments, indicates that the kinetics after  $n\pi^*$  and  $\pi\pi^*$  excitation of the (Z,Z)-isomer are identical after 1 ps and result in a main isomerization pathway with ~200 fs lifetime, followed predominantly by cooling dynamics in the ground state. The observed excited state dynamics is fully consistent with the computationally identified relaxation pathways for both  $n\pi^*$  and  $\pi\pi^*$  excitation of (Z,Z)-CBA. Furthermore, the computations provide a detailed insight into the molecular mechanism of the  $Z \rightarrow E$ photoisomerization reaction in CBA. The results contribute to the ongoing discussion of the photoisomerization pathways of ABs. They are also particularly relevant for the design of optically active macrocycles, where tailoring of the geometric distortions and constraints within the molecule can be exploited to control the functional properties (volume, rigidity and elasticity) of macrocyclic nanostructures. From a mechanochemical perspective, the results are an important step towards triggering transformations in the environment by mechanical forces exerted by photoswitching of macrocycles.

#### ASSOCIATED CONTENT

#### **Supporting Information**.

The following files are available free of charge.

Supporting information containing detailed description of the experimental and theoretical

methods and supporting figures. (file type, PDF)

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# ACKNOWLEDGMENT

C.S. and J.W. acknowledge funding by the DFG (WA 1850/4-2). C.Y. acknowledges the Chinese Science Council for his PhD fellowship and the support by the Heidelberg Graduate

School "Mathematical and Computational Methods for the Sciences" (GSC 2020).

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