



# **Fundamental Theoretical Investigations of the Photocatalytic deNO<sub>x</sub> Reactions over TiO<sub>2</sub> Surfaces**

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*by*

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## DECLARATION

I, Pu Guo, declare that the thesis entitled “ Fundamental Theoretical Investigations of the Photocatalytic deNO<sub>x</sub> Reactions over TiO<sub>2</sub> Surfaces” submitted by me, for the award of the degree of *Doctor of physics* to University of Bremen is a record of bonafide work carried out by me.

I further declare that the work reported in this thesis has not been submitted and will not be submitted, either in part or in full, for the award of any other degree or diploma in this institute or any other institute or university.

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This is to certify that the thesis entitled “Fundamental Theoretical Investigations of the Photocatalytic deNO<sub>x</sub> Reactions over TiO<sub>2</sub> Surfaces” submitted by Pu Guo, Bremen Center for Computational Material Science, University of Bremen, for the award of the degree of *Doctor of physics*, is a record of bonafide work carried out by her under my supervision.

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**(Prof. Thomas Frauenheim)**

## ABSTRACT

Emissions of nitrogen oxides ( $\text{NO}_x$ ) have a serious impact on our environment and human health. Photocatalytic denitrification ( $\text{deNO}_x$ ) on titania attracted much attention due to its low-cost and nonpolluting nature, but in reality undesired nitrite and nitrate were produced, instead of harmless  $\text{N}_2$ . Establishing the global activity and selectivity trend among active sites is an important base to explore and improve the  $\text{deNO}_x$  processes. Herein, I have investigated the reaction mechanisms of  $\text{NO}_2$  and  $\text{H}_2\text{O}$  on anatase  $\text{TiO}_2(101)$  and the direct decomposition of  $\text{NO}$  on various  $\text{TiO}_2$  surfaces. A polaron-corrected GGA functional (GGA + Lany-Zunger) was applied to improve the description of electronic states in photo-assisted process and a reaction phase diagram (RPD) was applied to understand the (quasi) activity trend over different active sites. In the  $\text{deNO}_2$  process, it was found that the perfect surface exhibits high activity while the activity on defective surfaces is limited by the sluggish recombinative desorption. A photogenerated hole can weaken the  $\text{OH}^*$  adsorption energies and circumvent the scaling relation of the dark reaction, ultimately significantly increasing the denitrification activity. In the  $\text{deNO}$  process, it was found that without illumination  $\text{N}_2$  production is inactive over various  $\text{TiO}_2$  surfaces/sites, and photo-generated holes can break the scaling relation of the dark condition by weakening  $\text{O}_2^*$  adsorption, leading to a significant increase in  $\text{deNO}$  activity of defective titania surfaces. The low  $\text{N}_2$  selectivity in experiments was further investigated by microkinetic modeling. It was found that the weak adsorption of  $\text{N}_2\text{O}$  led to the low  $\text{N}_2$  selectivity on titania surfaces. On the contrary, the  $\text{N}_2$  selectivity of Ti-modified zeolite was enhanced due to stronger  $\text{N}_2\text{O}^*$  adsorption. I show here that the reaction phase diagram (RPD) analysis can clearly establish a global picture of reaction activity and selectivity over various catalytic sites. The reactivity can be tuned by illumination-induced localized charge. Combined with microkinetic modeling, it can effectively determine the kinetic limits. This study provides insights for improving the design of photocatalysts.

**Keywords:** *Photocatalysis,  $\text{deNO}_x$ , Reaction phase diagram (RPD), GGA-LZ (Lany-Zunger) Functional, Microkinetic modeling.*

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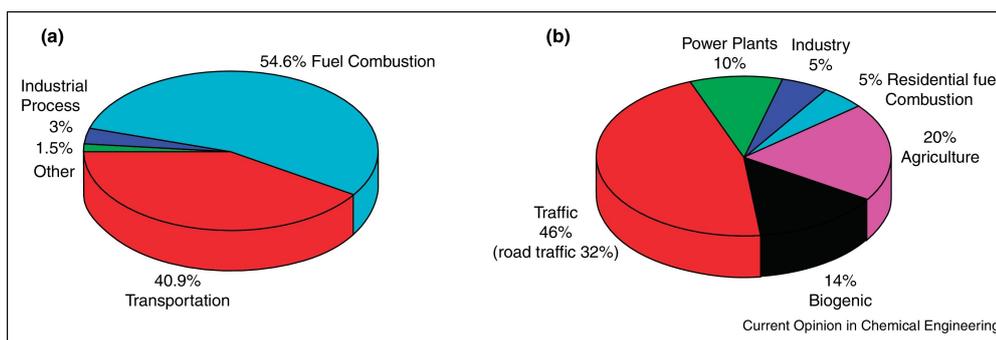
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# 1 | Introduction

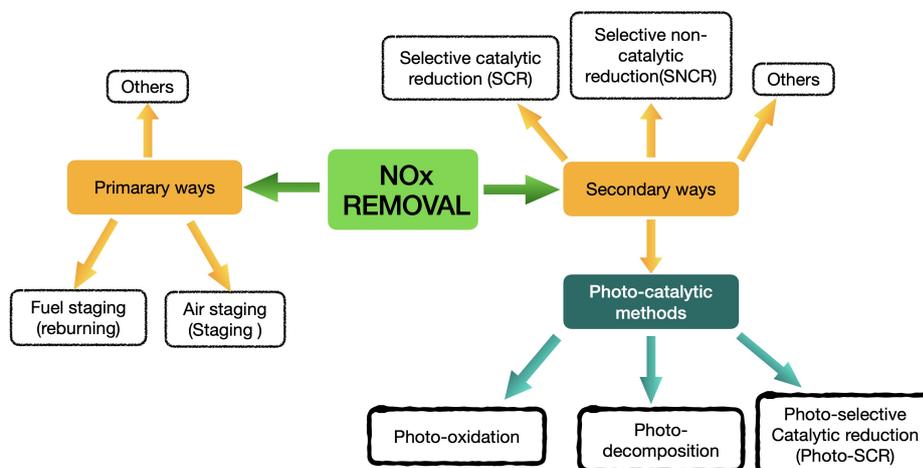


**Fig. 1.1** Nitrogen oxide emissions by source category in (a) the United States and (b) European Countries[1].

$\text{NO}_x$ , primarily  $\text{NO}$  and  $\text{NO}_2$ , are poisonous gases and have significant adverse effects on the environment and human beings. They mainly emitted into atmosphere through the combustion of fossil fuels as shown in **Fig. 1.1**.  $\text{NO}$  can be oxidized into  $\text{NO}_2$  in the air during the combustion processes and  $\text{NO}_x$  can also lead to acid rain due to its low solubility in water[1]. In addition,  $\text{NO}_x$  also affects the indoor areas. It has been reported that  $\text{NO}_2$  levels were 3-5 times higher in residential areas than rural area because of the cooking combustion[2].

With the increased interest toward removal of  $\text{NO}_x$  ( $\text{deNO}_x$ ), various methods had been applied in  $\text{deNO}_x$  process as shown in **Fig. 1.2**. The primary  $\text{NO}_x$  removal methods are carried out in the combustion zone without an additional reactor.  $\text{NO}_x$  production can be reduced by adjusting the combustion parameters, such as decreasing the temperature and the concentration of oxygen in the flame[2]. The prevailing secondary methods are selective catalytic reduction(SCR) and selective non-catalytic reduction(SNCR)[9]. Methods like  $\text{NO}$  absorption, phase separation[10],  $\text{NO}$  storage and reduction are all part of the other methods[2]. The main shortcoming of the secondary methods is that they must use additional equipment. Generally, the secondary methods are often more efficient in denitrification than primary methods[10, 2]. Photocatalysis is an attractive and promising technology for the removal of toxic and hazardous compounds[2]. There are three photocatalytic methods for photocatalytic denitrification: photo oxidation, photo selective catalytic reduction(photo-SCR), and photo

decomposition. Photo decomposition is an ideal process that does not take into account the presence of other compounds[2]. Therefore,  $\text{NO}_x$  can be oxidized to nitric acid in the presence of water or reduced to harmless  $\text{N}_2$  in the photocatalytic de $\text{NO}_x$  process.

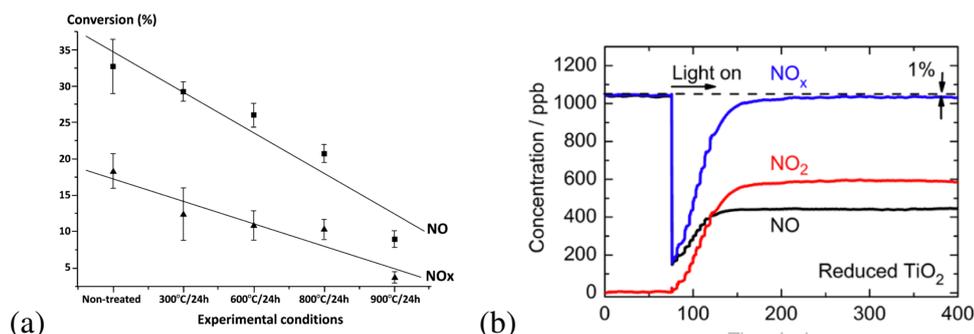


**Fig. 1.2** The methods of  $\text{NO}_x$  removal[2].

In the past decades,  $\text{TiO}_2$  photocatalysis has been extensively researched, and the process has been commercialized in several fields[11, 12, 13, 14, 15, 16, 17]. Toma et al. suggested that  $\text{NO}_x$  conversions depend on the crystalline phase of  $\text{TiO}_2$ [18]. As shown in **Fig. 1.3a**, the conversions of  $\text{NO}$  and  $\text{NO}_x$  in nontreated Degussa P25  $\text{TiO}_2$  were about 33% and 18%, respectively.  $\text{NO}$  and  $\text{NO}_x$  conversions decrease as the temperature increase. The comparison of de $\text{NO}_x$  mechanism between (anatase and rutile) titania surfaces can also provide useful guidance for catalyst design in the future. Thus, it is necessary to investigate the underlying effect of diverse crystalline phase of  $\text{TiO}_2$  on photocatalytic de $\text{NO}_x$  reactions.

Overall, the photocatalytic efficiency of  $\text{TiO}_2$  is not high. Enormous efforts had been made to improve the performance of  $\text{TiO}_2$ , for example, by metal or metal oxide deposition, elemental doping and defect formation[4]. It's reported that nanoscale anatase  $\text{TiO}_2$  and iron doped  $\text{TiO}_2$  have shown higher  $\text{NO}$  removal capacities[4, 19]. As shown in **Fig. 1.3b**, when the light was turned on the reaction proceeded immediately, implying that the UV-illumination light is indispensable for the photocatalytic de $\text{NO}_x$  reaction on reduced  $\text{TiO}_2$ [4]. Additionally, the de $\text{NO}_x$  reaction mechanisms has been widely discussed, but it is still controversial. The details is presented in the main text. In this work, I attempt to clarify the chemical reactions mechanisms of de $\text{NO}_x$  process and figure out the limits of it on diverse  $\text{TiO}_2$  surfaces.

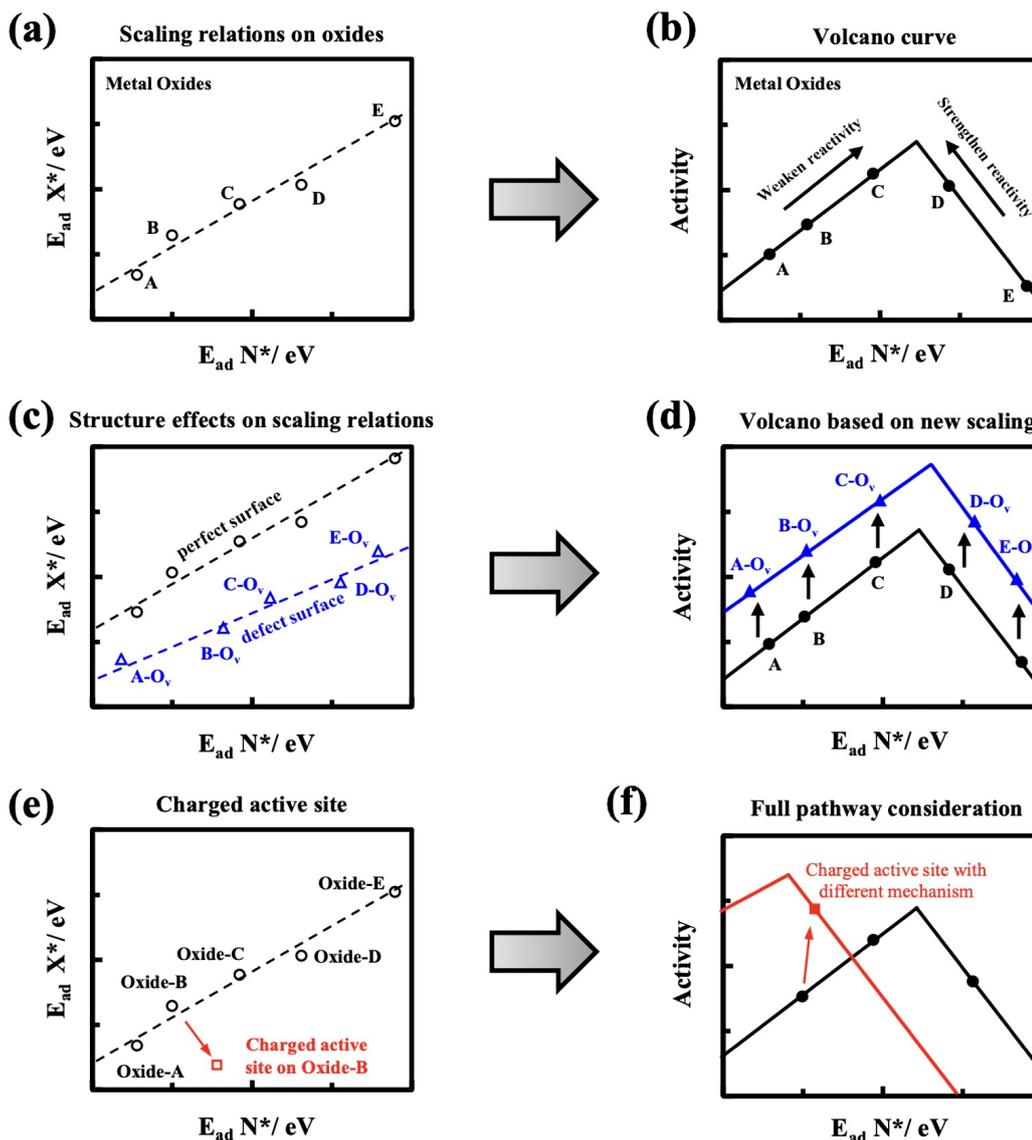
Computational catalyst design, based on electronic structure calculations, is becoming an important approach to provide guidance for the development of new materials with high activity and specific selectivity[5, 20, 21]. It is time consuming and hard to cover all possible reaction pathways if we study it in the traditional way. Therefore, it's



**Fig. 1.3** (a) Influence of the  $\text{TiO}_2$  annealing temperature on the photocatalytic conversion of nitrogen oxides, showing a negative correlation between the annealing temperature and the  $\text{TiO}_2$  photocatalytic activity. Reprinted with kind permission from Springer Science and Business Media[3]. (b) Photocatalytic degradation of NO in air by reduced  $\text{TiO}_2$  under UV light irradiation[4].

necessary to introduce an effective scheme to study a full set of reaction pathways. It was found that there are several linear correlations between adsorption energies of reactive adsorbates (e.g.  $C^*$  vs  $\text{CH}_x^*$ , described as:  $E_{ad} \text{CH}_x^* = a \times E_{ad} C^* + b$ ) (called scaling relations) over a series of catalysts (**Fig. 1.4a**) [5, 20, 22, 23, 24, 25, 26, 27], which can serve as the basis for reducing the dimensionality of complex reactions. One or two key adsorption energies can be chosen as descriptors to predict the activity trend. The well-known Sabatier principle can be quantitatively described through a volcano plot (**Fig. 1.4**) of rates calculated theoretically [21, 28, 29, 30, 31, 21, 32, 33, 34]. The term “volcano plot” comes from the shape of the curve describing the activity as a function of the adsorption energy, which was chosen as descriptor. The activity of the catalysts/sites in the left leg of a volcano plot (such as A, B and C in **Fig. 1.4a**) is limited by too strong reactivity, associated with the elementary processes of recombinative reactions, diffusion, and desorption. In contrast, the activity of catalysts/sites in the right leg (such as D and E in **Fig. 1.4b**) is limited by weak reactivity, resulting in a low coverage of reacting species and high barriers of dissociation [5, 20].

The scaling relation can be different due to the modification of local surface, like an oxygen vacancy, as shown in **Fig. 1.4c**. As a consequence, a new volcano plot (blue line in **Fig. 1.4d**) and reaction phase diagram can be achieved with improved activity. Oxygen defects of  $\text{TiO}_2$  have been the focus of interest due to charges trapping and the formation of small polarons [35, 36, 37, 38]. Charge-trapping centers on defective surfaces are usually considered as favored sites for adsorption and chemical reactions [39, 40, 41]. Note that oxygen vacancies are not typically present on the anatase  $\text{TiO}_2(101)$  surface [40], but can be created by photo beam under illumination condition [40, 42].



**Fig. 1.4** Strategy for computational investigation of catalysis. (a) Schematic diagram of linear correlations between adsorption energies (scaling relation) over oxide surfaces.  $N^*$  refers to the adsorbed N, which is the descriptor applied in the project,  $X^*$  refers to other adsorbates. (b) Traditional activity volcano plot for oxides. (c) New linear correlations between adsorption energies caused by local effects (e.g. oxygen vacancy). (d) The improved volcano plot based on new scaling relations over local effects. (e) Breaking the scaling relation on a charged system in photocatalysis. (f) Activity trend with the consideration of different reaction mechanisms[5].

In addition, the activity of  $\text{deNO}_x$  may be significantly different without and with illumination due to the effect of localized charge on the surface induced by the photo excitation, which may break the scaling relation of adsorption energies(Fig. 1.4a, e) and the volcano plot and shift the volcano plot to a higher activity.(Fig. 1.4f) The photogenerated charge may also give rise to a favored reaction pathway, which is different from the one under dark conditions[5, 7]. Therefore, unveiling the mechanism of pho-

tochemical deNO<sub>x</sub> on TiO<sub>2</sub> surfaces is very important for improving catalyst design. In such a study, a full diversity of reaction pathways should be considered, but so far, only some hypothetical and simplified reaction pathways have been used to study photocatalysis over oxides[43]. The varying reaction mechanisms and pathways assumed in the literature have led to great uncertainties in the calculated activity for a new material. In principle, it is possible to develop a scheme with full consideration of all reaction pathways; however, the computational cost of the construction of extremely complex reaction networks is becoming a major challenge. Herein, I applied our recently developed new algorithm[5], which can efficiently consider the full set of pathways. Furthermore, this approach shows the basic influences that limit the efficiency and activity of the catalysts[5, 20].

The scaling reactions and volcano plots are based on calculated adsorption and reaction energies, as well as energy barriers, and the predictive power is critically dependent on the accuracy of these calculations. The results can be quite inaccurate if standard local or semilocal implementations of density functional theory (DFT) are used (because those have a strong tendency to charge delocalization)[44] or if the artifacts caused by the repetition of the charge in periodic models of charge-assisted processes (like photocatalysis) is not corrected[45]. Therefore, it is quite important to have a computational strategy based on:

- 1). An effective scheme of studying a full set of reaction pathways, for which our group have developed a new algorithm[5].
- 2). Accurate methods for calculating the energetics beyond standard density functional theory and charge corrections which our group have recently developed[46, 47, 48].

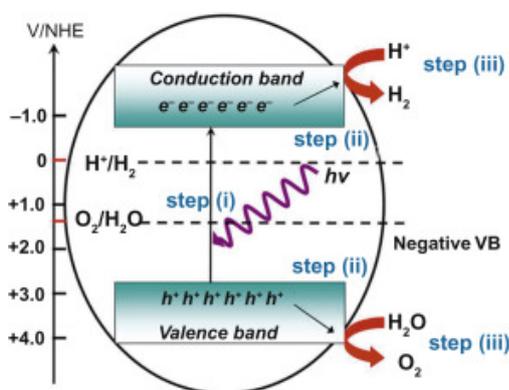
In this work, I used these aforementioned tools to address the highly important practical problem of increasing the activity of photocatalytic denoxification over oxides. I took into account a complete set of possible reaction pathways using a cutting-edge technology. Based on that, the calculated data obtained in a computational framework, which is capable to describe a huge number of charge-assisted reactions on oxide surfaces for the first time without the localization error of (as yet) standard electronic structure calculations, and completely free of the artifacts caused by imposing periodic boundary conditions on a charged system.

The thesis is organized as follows. I firstly introduce the fundamental principles of photocatalysis and the different crystalline phase of TiO<sub>2</sub> in Chapter 2, the computational methods and data analysis tools are discussed in Chapter 3, the computational parameters and models used to calculate the adsorption energy, to establish the reaction network, and to search for optimal reaction pathways are described in Chapter 4. The activity and selectivity maps of deNO<sub>2</sub> under dark and illumination conditions are given in Chapter 5. The activity and selectivity maps of deNO over diverse TiO<sub>2</sub> surfaces under dark and illumination conditions are given in Chapter 6. Microkinetic modelings of

photocatalytic deNO over titania surface and modified zeolite are given in Chapter 7. Finally, I present my conclusions and the insights obtained in this work as well as the implications for the photocatalysis field in general.

## 2 | Photocatalysis on titania

### 2.1 Basic principles in photocatalysis



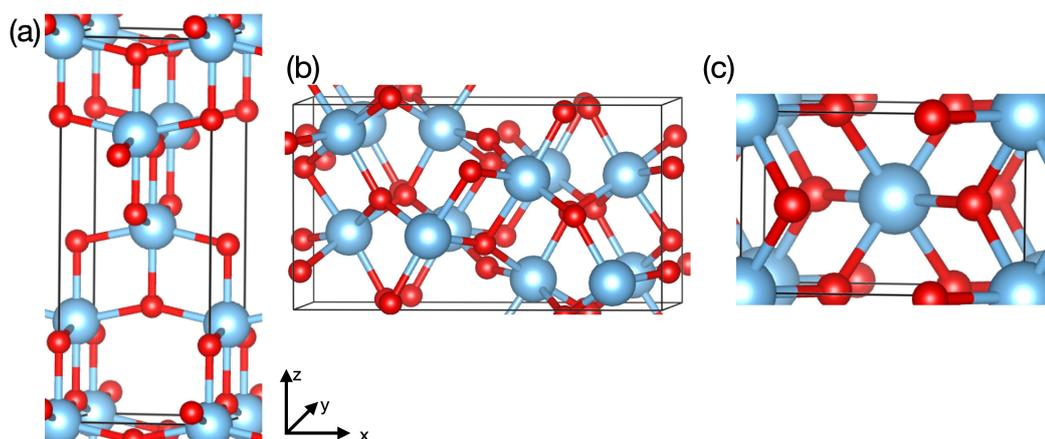
**Fig. 2.1** The schematic mechanism of photocatalytic water splitting. Step (i): excitation, step (ii): charge migration, step (iii): surface catalysis[6].

Photocatalysis, like photoassisted chemical reactions on a solid surface, relies on charge carriers generated by photoexcitation[6, 7]. The semiconductor will be excited as the energy of photons is equal to or larger than its band gap, leading to the generation of free moving electron-hole pair in the conduction band and valence band of the semiconductor, respectively. Theoretically, the more photons semiconductors can absorb, the higher the achieved efficiency of the semiconductor-based system[6]. The next step is defined as the “charge migration” step. The photogenerated electron and hole pairs migrate to the surface of the semiconductor. Note that the recombination of the electron and hole pairs should be inhibited in this step. Because the charge recombination significantly decreases the amount of the available active charges for further photocatalytic reactions. Therefore, it is crucial to propose strategies that can promote charge separation and transportation. This can be generally realized by p-n junction, loading cocatalysts and so on.

The last step is defined as the "surface catalysis" step. The electron can facilitate chemical reduction reactions and the hole can facilitate chemical oxidation reactions on the surface. A schematic illustration of photocatalytic water splitting is depicted in **Fig 2.1**. After photoexcitation, the photogenerated electron and hole can be used to

reduce  $H^+$  and oxidize  $H_2O$ , respectively. In principle, the bottom of the conduction band(CBM) must be higher than the  $H^+/H_2$  reduction potential (0 V vs. NHE at pH = 0) and the top of the valence band(VBM) must be lower than the  $O_2/H_2O$  oxidation potential (1.23 V vs. NHE at pH = 0). Therefore, a semiconductor with a band gap greater than the thermodynamic requirement of 1.23 eV and straddling the redox potential of  $H_2O$  would be an appropriate candidate for water splitting[6].

## 2.2 Introduction to titania



**Fig. 2.2** Crystal structure of anatase (a), brookite(b) and rutile(c)

Titanium dioxide (TiO<sub>2</sub>) is a low-cost, non-toxic, environmentally-friendly and simple-to-handle material. It has been widely investigated in photocatalysis since it was reported that TiO<sub>2</sub> could successfully decompose H<sub>2</sub>O with irradiation by Fujishima and Honda in 1972[49]. TiO<sub>2</sub> has three major polymorphs: anatase, rutile, and brookite (**Fig. 2.2**) with band gap of 3.2 eV, 3.0 eV and 3.4 eV, respectively[50]. In all these TiO<sub>2</sub> phases, the Ti atom is six-fold coordinated and is surrounded by six O atoms, forming a octahedron, TiO<sub>6</sub>[6]. The O atom is three-fold coordinated. The conduction band of TiO<sub>2</sub> mainly consists of Ti<sub>3d</sub> orbitals and the valence band consists of O<sub>2p</sub> orbitals. Rutile is the thermodynamically most stable phase, the other two forms of TiO<sub>2</sub> can convert to rutile when the heating temperature exceeds 600°C[18]. It is believed that anatase has a higher chemical activity than rutile[51]. Despite the fact that anatase is frequently noted as the most photo active polymorph, numerous studies report that rutile has comparable photo activities than anatase[3, 52]. Most commercial products of TiO<sub>2</sub> are mixtures of rutile and anatase, like Degussa P25 TiO<sub>2</sub>, which is mixed by 25% rutile and 75% anatase. The mixed phase has been widely used in laboratory studies because it usually performs better than any of the individual phase[3, 53]. In principle, it is hard to compare the performance of different polymorphs of TiO<sub>2</sub> under illumination conditions since many factors can affect the photo activity of TiO<sub>2</sub>, such

as the ratio of morphology, surface structure, and surface chemistry for different  $\text{TiO}_2$  samples as well as the target reactions[3]. At present, it cannot be concluded that the activity of anatase and rutile are essentially different. It would be quite helpful in the choice of experiments if the activity and selectivity of photocatalytic denitrification on various titania surfaces was understood.

The most stable facet of anatase  $\text{TiO}_2$  is the (101) surface that contributes around 94% of the exposed crystal surface[54]. The most stable facet of rutile  $\text{TiO}_2$  is (110) surface[55]. As mentioned before, oxygen vacancies are regarded as important adsorption sites[56], and active centers for reactions[4, 57], like water splitting at an oxygen vacancy[55, 58]. In principle, removing a bridging oxygen atom from  $\text{TiO}_2$  surfaces leads to the formation of two reduced  $\text{Ti}^{3+}$ [55], which make contributions to the donor levels in the electronic structure of  $\text{TiO}_2$ [59, 60]. Thus, much attention were paid to the development of synthesis of  $\text{TiO}_2$  with oxygen vacancies and exploration of its unique properties for photocatalytic applications. Thus, I focused primarily on defective and perfect anatase(101) and rutile(110) surfaces in this work.

## 3 | Computational methods

### 3.1 Density functional theory

Electronic structure calculations are extremely important in computational material science since most of the properties of a system are determined by its electronic structure, i.e. the spatial and energy distribution of electrons[61]. In 1926, Erwin Schrödinger firstly published his now-famous wave equation, which appeared to be extremely promising because it contained all available information about a system[62]. The Schrödinger equation was further simplified by the Born-Oppenheimer(BO) approximation[63], which took advantage of the significant differences between the masses of nucleus and electrons, separating the motion of the fast moving electron from that of the slow heavy nucleus. However, the Schrödinger equation is still difficult to solve even under the BO approximation when the number of electrons is large, due to the complex interactions between electrons. The difficulty of solving the Schrödinger equation grows exponentially with the number of electrons. Notably, Paul Dirac in 1929 attributed this difficulty to the lack of mathematical theory in dealing with a large part of physics and the whole of chemistry[64]. Then, the Hartree-Fock(HF) approximation, was utilized to solve the complex many-electron wave function. The HF theory replaces the real electrons with fictitious non-interacting electrons (so called one-electrons), whose wave functions are determined in such a way that the Slater determinant[65], set up from them, mimics the real many-body wave function. In other words, it assumed that there is no direct interaction between electrons and these non-interacting electrons move in a mean field of the rest, without knowing where the others are, i.e., without correlation between their motions. The total wave function can be expressed as a Slater determinant of  $N$  spin-orbitals to fulfill the Pauli-exclusion principle[65, 66, 67]. The disadvantages of HF methods are that it didn't cover the electron correlation and the computational cost is proportional to the 4th power of the number of electrons. Thus, it was desired to find a new method, which has to be rigorous but able to minimize the computational complexity as much as possible. Then, the density functional theory(DFT) was born.

In 1964, Hohenberg and Kohn published a landmark paper[68] and the theorems proven in this report are the bedrock for all modern density functional theories. The Hohenberg-Kohn theorems concerns any system consisting of electrons moving under

the influence of an external potential ( $V_{\text{ext}}(\mathbf{r})$ )[68]. The first HK theorem demonstrates that the ground-state properties of a many-electron system are uniquely determined by an electron density that depends on only three spatial coordinates. Namely, the ground-state wave function and the all other properties of the system defined by an external potential( $V_{\text{ext}}$ ) are determined by the ground state density.

In second theorem, they defined a universally valid functional,  $F[\rho]$ :

$$F[\rho] = T[\rho] + E_{ee}[\rho] \quad (3.1)$$

$$E[\rho] = \int \rho(r)V_{\text{ext}}dr + F[\rho]. \quad (3.2)$$

This functional reaches its minimum value if and only if the input density is the true ground state density. All the theorems tell us is that there is a unique mapping between the ground state density  $\rho_0(\mathbf{r})$  and the ground state energy  $E_0$ .

Unfortunately, the exact construction of  $F[\rho]$  is unknown. In 1965, Kohn and Sham made a significant breakthrough and they introduced the concept of a non-interacting reference system built from a set of orbitals (i. e., one electron functions) such that the major part of the kinetic energy can be computed to good accuracy[69]. The kinetic energy can be written as:

$$T_s = -\frac{1}{2} \sum_i^N \langle \phi_i | \nabla^2 | \phi_i \rangle \quad (3.3)$$

,where  $\phi_i$  is called Kohn-Sham orbitals, or briefly KS orbitals. The energy functional of the true system can be written as:

$$\begin{aligned} E[\rho(\vec{r})] &= T_s[\rho] + J[\rho] + E_{\text{xc}}[\rho] + E_{\text{ext}}[\rho] \\ &= T_s[\rho] + J[\rho] + E_{\text{xc}}[\rho] + \int \rho(r)V_{\text{ext}}dr \end{aligned} \quad (3.4)$$

,where  $J[\rho]$  is Coulomb interaction and  $J[\rho] = \frac{1}{2} \int \int \frac{\rho(r)\rho(r')}{|r-r'|} dr dr'$ ,  $E_{\text{xc}}[\rho]$  is exchange correlation energy and  $E_{\text{xc}}[\rho] \equiv (T[\rho] - T_s[\rho]) + E_{ee}[\rho] - J[\rho]$ . The exchange-correlation energy is also the functional of density and it contains everything that is unknown and also the the difference between the true kinetic energy and that of Eq.(3.3). The Kohn-Sham energy expression is subjected to the variation principle with respect to independent variations in the orbitals:

$$\left[ -\frac{1}{2}\nabla^2 + V_{\text{ext}}(r) + V_{\text{coul}}(r) + V_{\text{xc}}(r) \right] \phi_i(r) = \epsilon_i \phi_i(r) \quad (3.5)$$

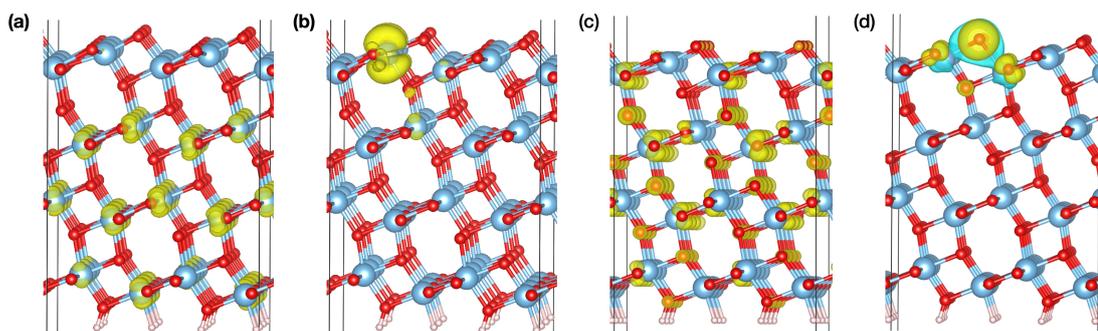
By this method, once we know the  $V_{\text{xc}}$ , the Kohn-Sham orbital can be obtained by solving the Kohn-Sham equation by the self-consistent field method. However, the

exchange-correlation energy can not be exactly calculated and approximations for exchange correlation functional must be used. Therefore, the main goal of modern density functional theory is to find better approximations for exchange correlation energy and the quality of the density functional approach depends entirely on the accuracy of the selected approximation to  $E_{XC}$ [70]. The main developed approximations are the local density approximation (LDA)[71], meta-GGA, general gradient approximation (GGA)[72], hybrid functional[73]("Perdew's ladder"). LDA can produce good structural properties but fails to produce binding energies due to overbinding[61]. GGA can provide reasonably accurate results within absolute average errors of about 5 kcal/mol for atomization energies[61] but it usually causes delocalization errors and underestimation of the band gap in semiconductors and insulators. Hybrid functionals incorporate some Hartree-Fock exchange (HF) and are semi-empirical in nature[74]. Hybrid functionals successfully provide the best performance in many applications but they are costly[61].

### 3.2 Lany-Zunger correction on GGA functional

As mentioned above, the inaccuracy of the standard local and semi-local approximations (LDA and GGA, respectively) leads to several problems, such as the underestimation of the band gap and delocalization of electron states. Thus, they are unable to describe small polarons (charge trapped in a strongly localized state by distortion of the pristine lattice)[75]. This problem can be avoided by using the hybrid exchange functional(HSE)[76, 36, 77], which can reproduce the band gap and fulfill the generalized Koopman's theorem (gKT)[78]. The gKT is a requisite for proper localization of states, predicting correct ionization energies[78, 79]. However, it is costly to simulate surface reactions with a HSE functional. GGA-LZ functional, a GGA functional with corrections introduced by Lany and Zunger, provides possibilities to the investigation of a complex reaction network. The band gap is corrected by using non-local external potentials[78], and occupation-dependent localizing potentials make sure that the gKT is satisfied. However, it was found that the localization potentials, with parameters determined for the bulk, do not satisfy the gKT for states localized on the surface, due to changes in the electronic screening. Thus, for properly treating surfaces, Meisam and co-workers have extended the Lany-Zunger (LZ) method by introducing coordination-dependent parameters for the localizing potentials[46]. This method can correctly predict charge localization both in the bulk and on the surface, and provides a good approximation of the much more expensive HSE results.

As shown in **Figure 3.1a and 3.1c**, a pure PBE(Perdew-Burke-Ernzerhof) calculation leads to the complete delocalization of an excess electron and hole over the entire bulk of the slab, whereas PBE + LZ describes the small polaron state well. A localized



**Fig. 3.1** The effect of the LZ-correction. Distribution of an extra electron in a pristine  $\text{TiO}_2(101)$  slab, as obtained from GGA without (a) and with (b) the Lany-Zunger correction. Distribution of an extra hole in a pristine  $\text{TiO}_2(101)$  slab, as obtained from GGA without (c) and with (d) the Lany-Zunger correction. The two spin directions are distinguished by color. Cyan, red, and pink denote Ti, O, and H, respectively.

electron state around a single Ti atom with strongly relaxed environment is shown in **Figure 3.1b** and a localized hole state around bridging O atom with relaxed environment is shown in **Figure 3.1d**. These results are consistent with higher level (HSE06) calculations by us and others[75, 76, 80]. This is an important improvement and the basis for analyzing charge-assisted photocatalytic reactions. The PBE(Perdew-Burke-Ernzerhof) + LZ calculation also reproduces the results for the neutral vacancy, as obtained by HSE06[81] and by screened exact exchange functional (sX) calculations[82]. It is obvious, that this has a profound effect on charge assisted surface reactions. Note that the delocalization tendency of plain GGA does affect the energetics of reactions in the dark too, if they are involving defects or adsorbates which give rise to levels in the band gap of  $\text{TiO}_2$ . HSE calculations in modelling heterogeneous catalysis are scarce and limited in scope. I applied GGA + LZ here for large scale calculations, with occasional checks using HSE.

### 3.3 Modeling charge assisted surface reactions

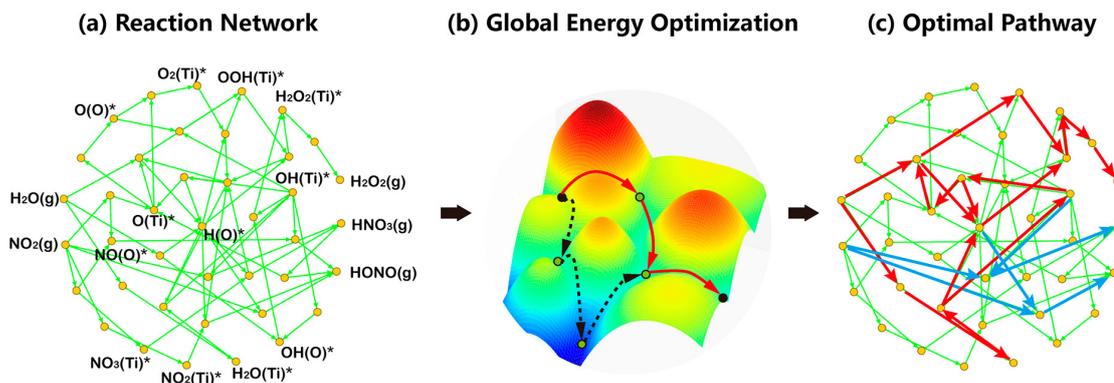
Herein, I used static electronic structure calculations in a model with a net charge to simulate the oxidation and reduction reactions, respectively. The net charge induces severe problems. For that reason, general computational studies of photocatalysis usually do not consider the effect of charge on the surface at all[80, 83, 84, 85], except when either a species with opposite charge is assumed in the gas phase[82] or a non-periodic model is used[86].

Adding a localized charge to a periodic model leads to an infinite Coulomb-energy, due to the electrostatic interaction between the periodic images[87]. To avoid that, in all electronic structure codes used for modelling periodic systems with a net charge, an evenly distributed (so called jellium) countercharge is added, to keep the repeated

supercell electrically neutral. Considering the different localization of the charge and countercharge, however, the repeated charges cause a spurious electrostatic potential, which affects the energy of the calculated one-electron states and shifts the total energy. For example, the formation energy of a charged Cl vacancy diverges with vacuum size and lateral size[87]. A posteriori method based on classical electrostatics was applied to correct bulk calculations with a charge[79, 88, 89, 90], but the correct treatment of a charge in the proximity of an interface is problematic because of the dielectric profile[87]. Our group recently have developed two procedures for charge corrections: SLABCC and SCPC. SLABCC calculates an a posteriori energy correction for charged slab models under 3D periodic boundary condition (PBC) based on the method of Komsa and Pasquarello[87], as implemented in the SLABCC code[47]. SCPC involves a potential correction, which automatically corrects both the total and the one-electron energies, as well as the wave functions, self-consistently and can be applied both in the bulk and for slabs. It has been implemented into the QE (Quantum Espresso) and VASP (Vienna ab initio simulation package) codes, two of the most commonly used electronic structure packages[48].

In this work, I employed these two charge correction methods and it was found that the new self-consistent potential correction (SCPC) method[48] have total energy corrections very close to the SLABCC results as shown in **Table A.1**.

### 3.4 Reaction phase graph method



**Fig. 3.2** (a) The complete deNO<sub>2</sub> reaction network with stable intermediates marked and transition states unmarked. (b) The algorithm of determining the pathways from global energy optimization, as described in Eq. 3.9. (c) The optimal reaction pathway for deNO<sub>2</sub> over the perfect anatase TiO<sub>2</sub>(101) surface (blue, NO<sub>2</sub> evolutionary paths; red, H<sub>2</sub>O evolutionary paths)[7].

To construct a full reaction network(**Figure 3.2**), the  $(m \times n)$  stoichiometric matrix,  $S$ , describing all elementary steps and species, was established. The rows and columns in  $S$  (Eq. 3.7) correspond to the substances( $m$ ) and the elementary steps( $n$ ) of the con-

sidered reactions, respectively. An elementary step  $j$  is described by the chemical vector,

$$R_j = (v_{1j}, v_{2j}, v_{3j}, \dots, v_{mj})^T \quad (3.6)$$

where  $v_j$  represents the stoichiometric number of the involved substances (reactants, intermediates, and products) in the relevant steps, and  $m$  represents substances in the reaction.  $T$  is the rank conversion of the matrix. Then, all possible pathways can be enumerated with free linear combinations of elementary step vector and the overall matrix  $S$  can be established as below[5]:

$$S = \begin{bmatrix} v_{1,1} & \dots & v_{1,m} \\ \vdots & \ddots & \vdots \\ v_{n,1} & \dots & v_{n,m} \end{bmatrix} \quad (3.7)$$

To speed up the computation, an accelerator strategy was developed. Two steps are performed prior to the enumeration of the combinations[5]:

- (1) The intermediates (substances  $m$ , excluding reactants and products), which are only considered in one elementary step, are eliminated. (Elementary steps elimination)
- (2) The elementary steps involved same intermediates  $B$  can be linked and pre-combined (e.g.  $A \rightarrow B$  and  $B \rightarrow C$ , where intermediate  $B$  only involves in these two reactions can be connected as  $A \rightarrow B \rightarrow C$ ). (Precombination). The full enumeration can then be done, following:

$$c_i = [\alpha_1, \alpha_2, \dots, \alpha_n], \alpha \in [1 \dots n] \quad (3.8)$$

where  $c_i$  and  $\alpha_n$  refer to a subset of elementary steps and to the chemical vector of the involved substances in the newly combined relevant steps, respectively.

A rank-based method were used to determine whether the individual combination leads to a possible pathway[5, 7]. The  $\Delta G_{RPD}$ -limiting steps are the reaction determining steps, which can be deduced from the reaction phase diagram (RPD) analysis as shown in **Figure 3.2b, c**. In principle, the optimal reaction pathway is the one with the minimal free energy change in the limiting step at a given site. Thus, the descriptor value ( $x$ ) could be identified using the following equation:

$$\Delta G_{RPD}(x) = \min_i(\max_j(\Delta G_{(i,j)}(x))) \quad (3.9)$$

where  $i$  denotes the possible reaction free energies, which were obtained using all forward and backward elementary steps. All possible reaction pathways were obtained by the free combination of all the considered elementary steps.  $j$  denotes the elementary step with the maximum  $\Delta G_{RPD}$  ( $\Delta G_{RPD}$ -limiting step) in each pathway, which was selected and compared with all other pathways to minimize the  $\Delta G_{RPD}$ -limiting ener-

gies for total considered reactions. For instance, the pathways A and B consist of (R1 + R2 + R3 + R4) and (R5 + R6 + R7 + R8), respectively. R2 and R5 are most unfavorable over all elementary steps in internal comparison for the pathways A and B, respectively, while R5 is more favorable with respect to R2. As a result, pathway B is defined as the optimized reaction channel and R5 is the  $\Delta G_{RPD}$ -limiting step. Herein, the  $\Delta G_{RPD}$ -limiting energies over all active sites were used to determine the (quasi) activity trend and mechanism variation under dark and illumination conditions[7].

### 3.5 Microkinetic simulations

Microkinetic modeling has been widely applied in complex reaction mechanisms in last few decades[43, 91, 92]. It plays an important role in understanding the activity and selectivity of catalysts and in catalyst design[43]. Because it is possible to quantify the importance of not only each elementary step, but also each intermediate and transition state in the reaction mechanism by a general concept degree of rate control(DRC)[93, 94]. The definition of DRC[93] is :

$$X_i = \left( \frac{\delta(\ln r)}{\delta(-G_i^0/RT)} \right) \quad (3.10)$$

or

$$X_i = \frac{\delta(\ln r)}{\delta g_i} \quad (3.11)$$

$$g_i = -G_i^0/RT \quad (3.12)$$

, where r is the net steady-state rate for consumption of one reactant.  $G_i^0$  is the standard-state Gibbs free energy of species i.

A microkinetic simulation can be performed by the following steps:

- Proposing a reaction mechanism.
- Calculating all the reaction energies and reaction energy barriers.
- Using CATKINAS package[95] to simulate the reaction with practical industrial or experimental catalytic conditions.
- Finally, the CATKINAS package outputs the results including activity, coverage, reversibility, sensitivity, energy profile and reaction flow diagrams.

### 3.6 Calculation of adsorption energy

In the deNO<sub>2</sub> process, the adsorption energies of the relevant intermediates on the solid surface, with respect to the gas phase energies of NO<sub>2</sub>, H<sub>2</sub>O<sub>2</sub> and H<sub>2</sub>O, were obtained from the following equation:

$$E_{\text{ads}} = E_{\text{tot}} - E_{\text{sur}} - (x\mathbf{E}_{\text{H}} + y\mathbf{E}_{\text{O}} + z\mathbf{E}_{\text{N}}) \quad (3.13)$$

, where  $x$ ,  $y$  and  $z$  denote the number of H, O and N atoms, respectively.

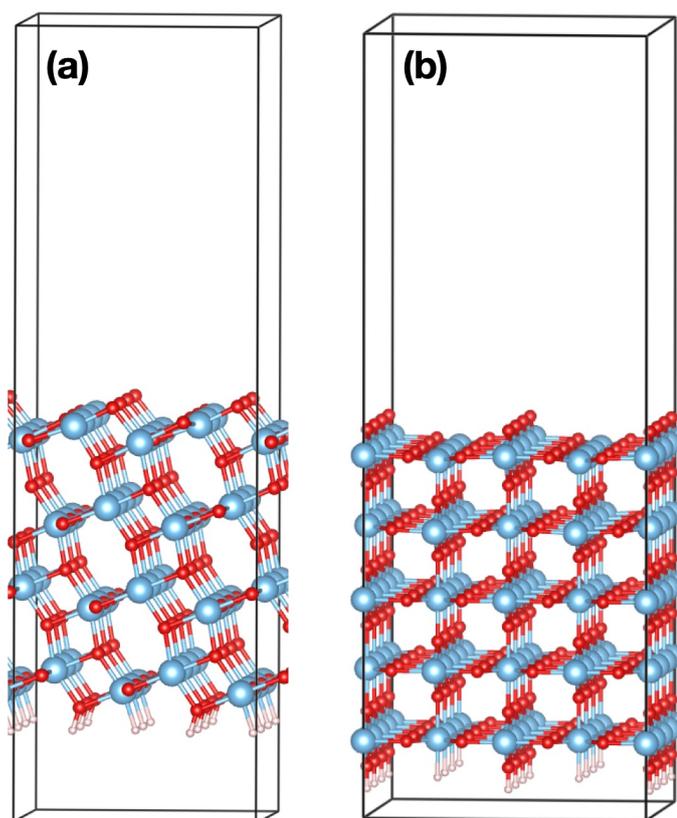
In the deNO process, the adsorption energies of relevant intermediates on the solid surface referred to the gas phase energies of NO and O<sub>2</sub>, which were obtained from the following equation:

$$E_{\text{ads}} = E_{\text{tot}} - E_{\text{sur}} - (xE_{\text{O}} + yE_{\text{N}}) \quad (3.14)$$

, where  $x$  and  $y$  denote the number of O and N atoms, respectively.

The thermal correction for Gibbs free energy( $G$ ) includes the zero point energy(ZPE) and Gibbs free energy difference between 0 K and T K, i.e., the correction for  $G(T) = \epsilon_{\text{ZPE}} + \Delta G_{0 \rightarrow T}$ . The thermal corrections were obtained using the methods implemented in VASPKIT program[96]. The temperature ( $T$ ) and pressure ( $P$ ) were set according to the relevant experimental conditions. In the deNO<sub>2</sub> process,  $T = 298\text{K}$ ,  $P = 202.65 \text{ nbar}$ ,  $9.12 \text{ nbar}$ ,  $2.53 \text{ nbar}$ ,  $9.12 \text{ nbar}$  and  $3.04 \text{ nbar}$ , for NO<sub>2</sub>, H<sub>2</sub>O, H<sub>2</sub>O<sub>2</sub>, HONO, and HNO<sub>3</sub> respectively[97]. In the deNO process,  $T = 493\text{K}$ ,  $P = 1.0\text{E-}6 \text{ atm}$  for NO[4]. The climbing image nudged elastic band (CI-NEB) method[98, 99, 100] and dimer method[101] were applied to search for transition states (TSs). The TSs were confirmed by frequency calculations to have only one imaginary frequency.

## 4 | Computational setup



**Fig. 4.1** (a) The relaxed structures of the anatase(101) (a), and rutile(110) (b) slabs. Cyan, red, and pink denote Ti, O, and H, respectively.

The lattice parameters  $a = 3.741 \text{ \AA}$  and  $c = 9.526 \text{ \AA}$  for anatase and  $a = 4.571 \text{ \AA}$  and  $c = 2.960 \text{ \AA}$  for rutile were used. The anatase (101) surface was modeled by a slab of four bi-layers with a  $3 \times 1$  surface supercell. The rutile (110) surface was modeled by a slab of five layers with a  $4 \times 2$  surface supercell. A vacuum layer of  $20 \text{ \AA}$  was applied. The two bottom layers were fixed and the five-fold coordinated Ti ( $\text{Ti}_{5c}$ ) and two-fold coordinated O ( $\text{O}_{2c}$ ) atoms were saturated by  $\text{H}^{1.33}$  and  $\text{H}^{0.66}$  pseudo-hydrogen atoms, respectively to avoid surface state. The optimized slab models are shown in **Figure 4.1**. All optimized adsorption geometries are shown in **Appendix B, C**. The valence electrons ( $10e^-$  for Ti and  $6e^-$  for O) were treated by plane

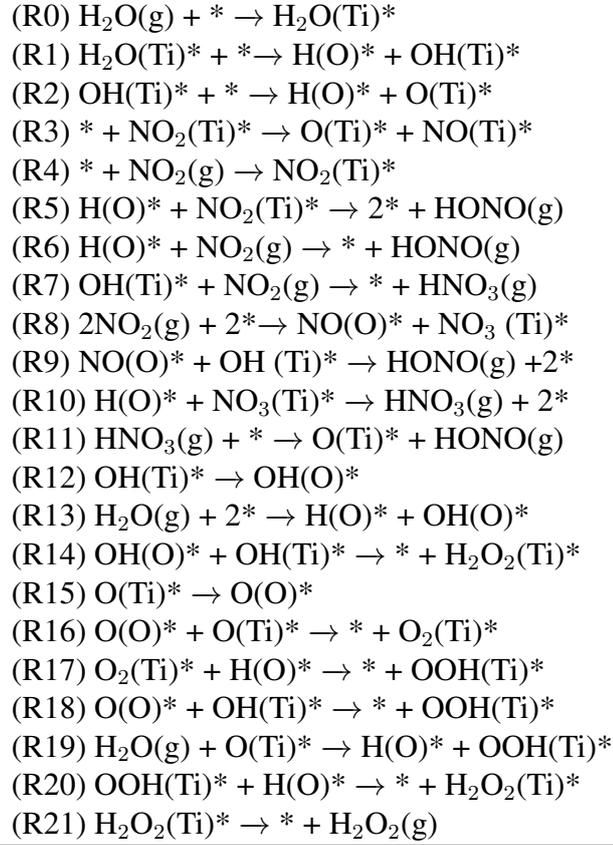
wave based GGA calculations, using the PBE exchange correlation functional[102]. The core-electrons were treated using the projector augmented wave method implemented in the VASP package[103, 104, 105]. A cutoff of 420 eV was used for the plane wave expansion and a cutoff of 840 eV was used for the charge density representation in the real space grid. GGA + U were applied in calculations of zeolite clusters with the U value for Ti 4.20 eV[106]. Lany-Zunger corrections were applied in calculations on the TiO<sub>2</sub> surface to reproduce the band gap and localize the polarons in an efficient way[78, 79]. SLABCC[47] and SCPC[48] methods were applied to correct the total energies in charged system as shown in **Appendix A**.

## 5 | Activity and reaction mechanism mapping of photocatalytic deNO<sub>2</sub> processes on anatase TiO<sub>2</sub>(101) surface

As we know that NO<sub>2</sub> is more toxic than NO, it heavily affects the overall deNO<sub>x</sub> processes[14]. Thus, NO<sub>2</sub> conversion on TiO<sub>2</sub> was focused on firstly. In experiments, various techniques have been applied to study NO<sub>2</sub> adsorption on TiO<sub>2</sub> particles and surfaces such as fourier transform infrared (FTIR) spectroscopy and XPS[12, 107]. It was reported that HONO and HNO<sub>3</sub> was produced on TiO<sub>2</sub> via the reaction of H<sub>2</sub>O and NO<sub>2</sub> under dark conditions[97, 107, 108, 109]. It might be also noted that HONO is an important source of reactive OH• radicals and key oxidant in the degradation of most air pollutant in the troposphere, which makes HONO formation reaction attach great attention[11]. The reaction mechanism of dark condition has been widely discussed in published papers but it's still controversial[110, 111, 112, 113, 114]. It was proposed that NO<sub>2</sub> disproportionation on TiO<sub>2</sub> has led to the formation of NO<sup>+</sup> and NO<sub>3</sub><sup>-</sup>, where the formed anion and cation can immediately react with H<sub>2</sub>O to produce HONO and HNO<sub>3</sub> in the subsequent steps[111, 110, 112, 113]. Haubrich et al have observed N<sub>2</sub>O<sub>4</sub> at higher NO<sub>2</sub> pressure by XPS[12]. Furthermore, they suggested HNO<sub>3</sub> is formed by physisorbed NO<sub>2</sub> rather than chemisorbed NO<sub>2</sub> and NO<sub>3</sub> species[12]. However, Lambert suggested that HONO was formed by direct hydrogenation of NO<sub>2</sub> under dark condition, where the hydrogen was generated from water dissociation[97]. It was detected that both oxygen atoms in the product of HONO can stem from NO<sub>2</sub> rather than water by isotope-labelling experiments. This suggested that the observed N<sub>2</sub>O<sub>4</sub> cannot be accounted for NO<sub>2</sub> hydrolysis under realistic pressures.

The photocatalytic mechanism is also disputed. Under UV illumination condition, it was reported by Lambert's group that H<sub>2</sub>O<sub>2</sub> is a minor product and they proposed H<sub>2</sub>O<sub>2</sub> was formed by the combination of OH radicals, arising from photocatalytic water splitting. Meanwhile, NO<sub>2</sub> captured the photogenerated electron and combined with H<sup>+</sup>

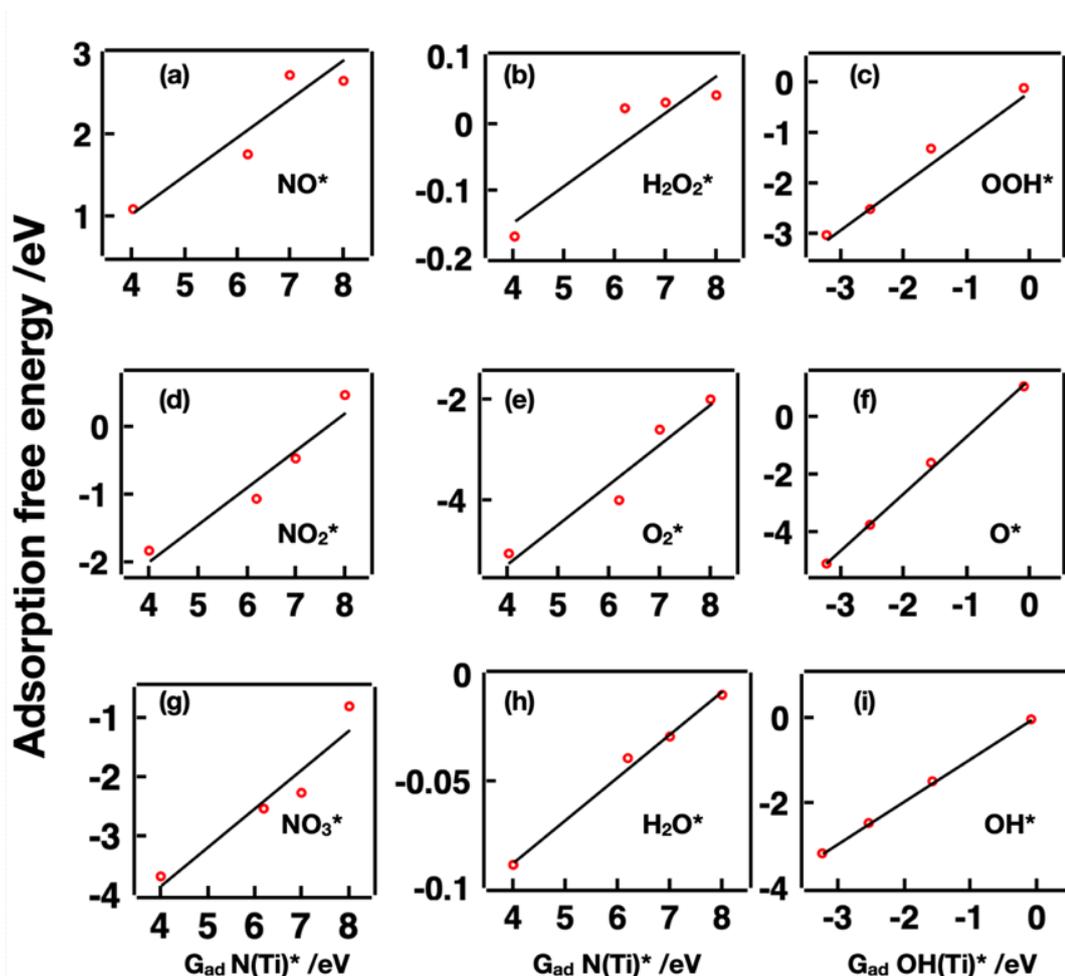
to produce HONO[97]. However, Bedjanian and co-workers reported the main product was HONO with illumination and they didn't detect H<sub>2</sub>O<sub>2</sub>[109]. Therefore, deNO<sub>2</sub> with the presence of water on TiO<sub>2</sub> surface is complex and not yet clearly understood and further studies are needed to understand them.



**Table 5.1** The considered elementary steps for the deNO<sub>2</sub> reactions on anatase TiO<sub>2</sub>(101) surface.

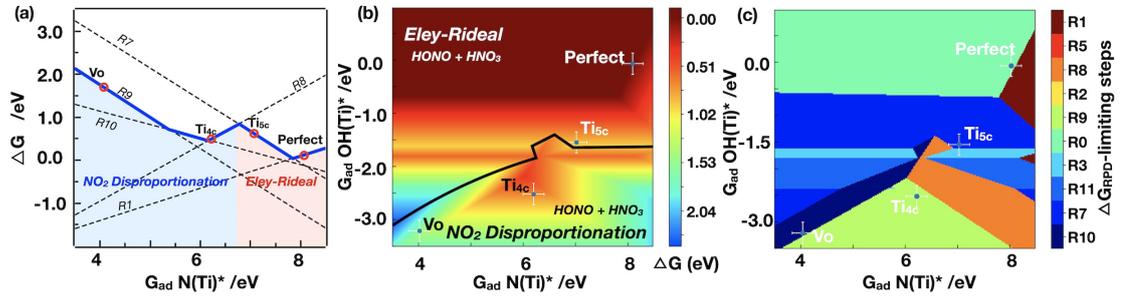
To simplified the reaction network, 22 possible elementary steps were considered in the deNO<sub>2</sub> reaction network as listed in **Table 5.1**. I obtained a total of 529 possible reaction pathways via the free combination of the all forward and backward elementary steps, and the complete reaction network is shown in **Fig. 3.2**. In the preliminary work, I also considered the combination of two species that are both adsorbed to Ti sites and other migration steps. However, due to the higher thermodynamic energy barriers, these elementary steps are not included in the optimal reaction pathways, so I have removed them in the following work. The adsorption energies of all important intermediates (H\*, O\*, NO<sub>2</sub>\*, OH\*, OOH\*, NO\*, NO<sub>3</sub>\*) and adsorbed molecules (H<sub>2</sub>O\*, H<sub>2</sub>O<sub>2</sub>\*) were calculated for different active sites on the surface of defective and perfect anatase TiO<sub>2</sub>(101). There are five considered active sites, such as oxygen vacancy(Vo), four-fold coordinated Ti (Ti<sub>4c</sub><sup>defect</sup>), five-fold coordinated Ti of defective surface (Ti<sub>5c</sub><sup>defect</sup>), five-fold coordinated Ti of perfect surface (Ti<sub>5c</sub><sup>perfect</sup>) and bridging oxygen(O<sub>br</sub>). To estimate the reaction free energies and improve the accuracy, two-dimensional fitting

was applied to establish the correlations between the adsorption free energies and the descriptors. The adsorption energies of  $\text{N}(\text{Ti})^*$  and  $\text{OH}(\text{Ti})^*$  were chosen as descriptors for the adsorption energies of all reaction species. As shown in **Fig. 5.1**, the descriptor for the adsorption free energy of  $\text{NO}^*$ ,  $\text{H}_2\text{O}_2^*$ ,  $\text{NO}_2^*$ ,  $\text{NO}_3^*$ ,  $\text{O}_2^*$  and  $\text{H}_2\text{O}^*$  was  $G_{\text{ad}} \text{N}(\text{Ti})^*$  and the descriptor for adsorption free energy of  $\text{OOH}^*$ ,  $\text{O}^*$ ,  $\text{OH}^*$  was  $G_{\text{ad}} \text{OH}(\text{Ti})^*$ . The scaling relations show good linearity. Note that the adsorption free energies of the bridging oxygen sites on perfect and defective surfaces of the considered species are considered to be the same, since the calculated energies were quite close. Based on the scaling relations, the free energy change of each elementary step can be described by descriptors.



**Fig. 5.1** Scaling relations of the adsorption free energies for (a)  $\text{NO}^*$ , (b)  $\text{H}_2\text{O}_2^*$ , (c)  $\text{OOH}^*$ , (d)  $\text{NO}_2^*$ , (e)  $\text{O}_2^*$ , (f)  $\text{O}^*$ , (g)  $\text{NO}_3^*$ , (h)  $\text{H}_2\text{O}^*$ , (i)  $\text{OH}^*$  of different active sites on the perfect and defective anatase  $\text{TiO}_2(101)$  surface. All the adsorption energies were calculated with reference to the energy of the gas molecules  $\text{NO}_2$ ,  $\text{H}_2\text{O}_2$  and  $\text{H}_2\text{O}$ . The adsorption free energies of  $\text{N}(\text{Ti})^*$  and  $\text{OH}(\text{Ti})^*$  were chosen as descriptors.

### 5.1 Activity and mechanism mapping of NO<sub>2</sub> and H<sub>2</sub>O without illumination on anatase TiO<sub>2</sub>(101) surface



**Fig. 5.2** (a) The reaction free energies of key elementary steps on all sites without illumination. Based on the scaling relationships between the adsorption free energies of adsorbates and descriptor ( $G_{\text{ad}} N(\text{Ti})^*$ ), the reaction free energies ( $\Delta G$ ) can be described by the descriptor ( $G_{\text{ad}} N(\text{Ti})^*$ ). The bold blue line denotes the 1D  $\Delta G_{\text{RPD}}$ -limiting steps for NO<sub>2</sub> + H<sub>2</sub>O reactions on different sites, which are determined with Eq. 3.9. “NO<sub>2</sub> disproportionation” and “Eley-Rideal” denote that the favored pathways follow NO<sub>2</sub> disproportionation and Eley-Rideal mechanisms, respectively. (b) The 2D (quasi) activity map without illumination is shown with two descriptors: adsorption free energies of N(Ti)\* and OH(Ti)\*, independently. The color bar denotes the limiting reaction free energies. (c) The 2D  $\Delta G_{\text{RPD}}$ -limiting steps for the NO<sub>2</sub> + H<sub>2</sub>O reactions on different sites without illumination. Vo denotes the oxygen vacancy; Ti<sub>4c</sub> and Ti<sub>5c</sub> denote the 4-fold and 5-fold coordinated Ti site on the defective surface, respectively. “Perfect” denotes the 5-fold coordinated Ti site on the perfect surface. The corresponding elementary steps are listed in **Table 5.1**

As shown in **Fig. 5.2a, b**, it was found that all sites were selective to HONO + HNO<sub>3</sub> production under dark conditions, which is in agreement with the experimental report [97], while the perfect surface exhibits the highest activity. Water dissociation on defective surfaces was thermodynamically favored, as shown in **Fig. 5.2a**. This is consistent with previous reports that water dissociation required a lower barrier on the reduced anatase TiO<sub>2</sub>(101) surface[39, 41]. On the defective surface, the deNO<sub>2</sub> activity on Ti<sub>5c</sub> and Ti<sub>4c</sub> is relatively better than on Vo with the  $\Delta G_{\text{RPD}}$ -limiting energies of 0.87 eV and 0.69 eV, respectively. As shown in **Fig. 5.2c**, the main  $\Delta G_{\text{RPD}}$ -limiting steps are:



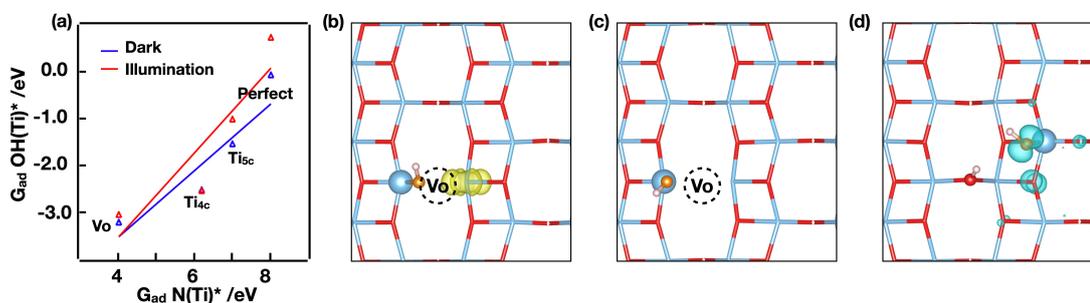
and



On the defective surface, the OH\* species has a too strong interaction with the defective surface sites, which hinders recombinative desorption. (Note that the labeling of the reactions follows that of **Table 5.1**.) Therefore, the defective surface, especially the Vo site, exhibits strong reactivity but low catalytic activity for deNO<sub>2</sub> reactions. On the perfect surface, the  $\Delta G_{\text{RPD}}$ -limiting step is water adsorption. In general, there

are multiple activity optima because of the diverse reaction pathways. The optimal reaction pathways to produce HONO and  $\text{HNO}_3$  were identical on the  $\text{Ti}_{5c}$  site of the defective surface and the perfect surface. HONO was formed by the combination of  $\text{H}^*$  and  $\text{NO}_2(\text{g})$ , while  $\text{HNO}_3$  was produced by  $\text{OH}(\text{Ti})^* + \text{NO}_2(\text{g}) \rightarrow \text{HNO}_3(\text{g}) + *$  via the Eley-Rideal mechanism. This is consistent with Habrich's observation on the rutile  $\text{TiO}_2(110)$  surface[115, 116]. Namely, the  $\text{HNO}_3$  production required weakly adsorbed  $\text{NO}_2$  species[115]. Note that the strong adsorption energy of  $\text{OH}(\text{Ti})^*$  inhibits the recombinative desorption (R9 and R7), which is adverse for the  $\text{deNO}_2$  activity. On  $\text{Ti}_{4c}$  and  $\text{Vo}$  sites, HONO and  $\text{HNO}_3$  were produced via the disproportionation of  $\text{NO}_2$  (R8), where the generated  $\text{NO}(\text{O})^*$  and  $\text{NO}_3(\text{Ti})^*$  species reacted with the dissociated water to produce HONO and  $\text{HNO}_3$ . Note that some sites are located at the boundary of multiple elementary steps and pathways (Fig. 5.2b, c), which indicates that those elementary steps and pathways exhibit similar reaction activities. In other words, it is possible to produce HONO and  $\text{HNO}_3$  via the disproportionation of  $\text{NO}_2$  on the  $\text{Ti}_{5c}$  site, which is also the  $\Delta G_{\text{RPD}}$ -limiting step with a reaction free energy of 1.01 eV.

## 5.2 Activity and mechanism mapping of $\text{NO}_2$ and $\text{H}_2\text{O}$ with illumination on anatase $\text{TiO}_2(101)$ surface

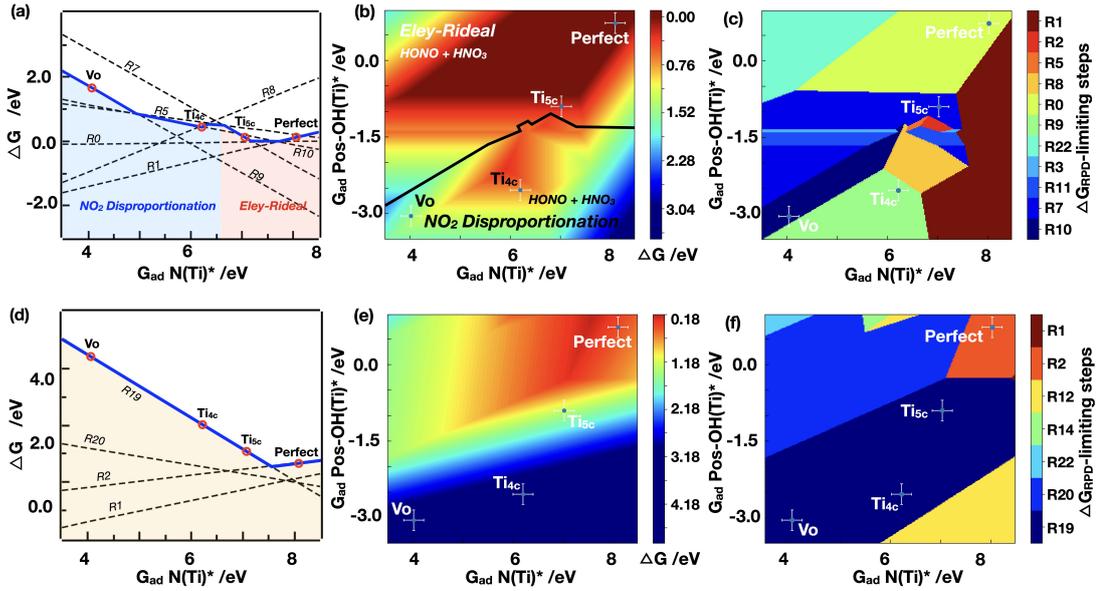


**Fig. 5.3** (a) The correlation between  $G_{\text{ad}} \text{N}(\text{Ti})^*$  and  $G_{\text{ad}} \text{OH}(\text{Ti})^*$  under dark and illumination conditions. Spin density of  $\text{OH}^*$  on (b) the neutral anatase  $\text{TiO}_2(101)$  slab with oxygen vacancy and (c) on the positively charged one. (d) Spin density of dissociated water adsorbed on perfect anatase  $\text{TiO}_2(101)$  slab with a positive charge. Yellow and turquoise indicate positive (spin up) and negative (spin down) values, respectively. Cyan and red denote Ti and O of  $\text{TiO}_2(101)$ , respectively. Orange and pink denote the O and H of  $\text{OH}^*$ , respectively.

The combination of elementary steps (possible pathways) is a pure mathematical solution, which can not be different for illumination conditions. However, the most favorable pathway is indeed possible upon switching to another. With our algorithm, all possible pathways are determined prior to energetic comparison. It is feasible to reanalyze the most favorable pathway under illumination condition. The photogenerated charges can influence a given reaction pathway in multiple times, namely, multiple

photon processes. As the limiting step in dark conditions is the most important, I studied the effect only on the limiting step by a single photogenerated charge. As OH(Ti)\* is a key species under dark conditions, I examined the G<sub>ad</sub> OH(Ti)\* in the presence of a photogenerated hole. At the Ti<sub>4c</sub> and Vo sites, G<sub>ad</sub> OH(Ti)\* does not change, as shown in **Fig. 5.3a**. However, it is significantly weakened on the Ti<sub>5c</sub> site of the perfect surface and on the Ti<sub>5c</sub> site of the defective surface, and this has a crucial role in enhancing deNO<sub>2</sub> activity. As can be seen from **Fig. 5.3c**, the photogenerated hole on the defective surface is trapped by the reduced Ti, against our expectation, rather than on the OH(Ti)\* (negative), so the G<sub>ad</sub> OH(Ti)\* remains intact. On the perfect surface, the hole is trapped by the OH(Ti)\* (negative), giving rise to an OH radical, as shown in **Fig. 5.3d**. Therefore, the illumination has no significant effect on the activity of the Vo and Ti<sub>4c</sub> sites, because the ΔG<sub>RPD</sub>-limiting energies can not be lowered. However, the activity on the perfect surface can be enhanced.

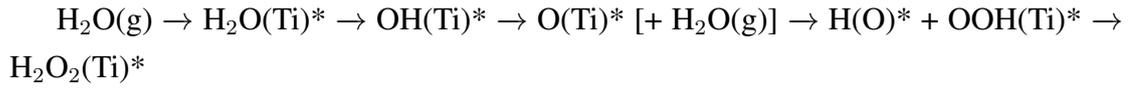
Assuming that water adsorption is not directly affected by illumination, its dissociation was considered as a stepwise reaction[117]: deprotonation first and then oxidation: OH(Ti)\* + hole → Pos-OH(Ti)\* (R22) while the photogenerated electron was trapped by the proton. The adsorption free energies of OH(Ti)\*, calculated on the positively charged surfaces (named as Pos-OH(Ti)\*), and the adsorption free energy of N(Ti)\* were chosen as descriptors. The reaction phase diagrams were established as shown in **Fig. 5.4a, b**, to analyze the deNO<sub>2</sub> activity in the photoassisted process. The perfect surface still exhibits the highest activity, and the limiting step is the same as under dark conditions. On the defective surface, the deNO<sub>2</sub> activity is enhanced on the Ti<sub>5c</sub> site and the ΔG<sub>RPD</sub>-limiting step has also changed, as shown in **Fig. 5.4c**. However, the Vo site remained inactive because of the too strong bonding interaction for Pos-OH(Ti)\* species, which is consistent with our hypothesis above. The main production pathways of HONO and HNO<sub>3</sub> on all sites was unaffected by illumination. At Ti<sub>5c</sub>, HONO and HNO<sub>3</sub> still prefer to be produced via H(O)\* + NO<sub>2</sub>(g) → \* + HONO(g) (R6) and OH(Ti)\* + NO<sub>2</sub>(g) → \* + HNO<sub>3</sub>(g) (R7) rather than via the NO<sub>2</sub> disproportionation (R8).



**Fig. 5.4** (a) The reaction free energies of key elementary steps on all sites in the deNO<sub>2</sub> reactions under illumination conditions. Based on the scaling relationships between the adsorption free energies of adsorbates and descriptor ( $G_{\text{ad}} N(\text{Ti})^*$ ), the reaction free energies ( $\Delta G$ ) can be described by the descriptor ( $G_{\text{ad}} N(\text{Ti})^*$ ). The bold blue line denotes the 1D  $\Delta G_{\text{RPD}}$ -limiting steps (determined by Eq. 3.9) for  $\text{NO}_2 + \text{H}_2\text{O}$  reactions on different sites. “NO<sub>2</sub> disproportionation” and “Eley-Rideal” denote that the reaction follows the NO<sub>2</sub> disproportionation and Eley-Rideal mechanisms, respectively. (b) The 2D (quasi) activity map with illumination is shown with two independent descriptors: adsorption free energies of N(Ti)\* and Pos-OH(Ti)\*. The Pos-OH(Ti)\* denotes the adsorption free energy of OH(Ti)\* in the positively charged system. The color bar denotes the limiting reaction free energies. (c) The 2D  $\Delta G_{\text{RPD}}$ -limiting steps for the deNO<sub>2</sub> reactions on different sites under illumination condition. (d) The reaction free energies of key elementary steps for  $2\text{NO}_2 + 2\text{H}_2\text{O} \rightarrow \text{H}_2\text{O}_2 + 2\text{HONO}$  over all sites under illumination condition. (e) The 2D (quasi) activity map of H<sub>2</sub>O<sub>2</sub> production with illumination is shown with two descriptors: adsorption free energies of N(Ti)\* and Pos-OH(Ti)\*, independently. The color bar denotes the limiting reaction free energies. (f) The  $\Delta G_{\text{RPD}}$ -limiting steps for  $2\text{NO}_2 + 2\text{H}_2\text{O} \rightarrow \text{H}_2\text{O}_2 + 2\text{HONO}$  over all sites with illumination. The corresponding elementary steps are listed in **Table 5.1**.

### 5.3 H<sub>2</sub>O<sub>2</sub> production during the photochemical reduction of NO<sub>2</sub> to HONO on anatase TiO<sub>2</sub>(101) surface with illumination

I studied the H<sub>2</sub>O<sub>2</sub> production separately. The 2D quasi activity map was established and is shown in **Fig. 5.4e**. Activity of the H<sub>2</sub>O<sub>2</sub> production is lower than that of HNO<sub>3</sub> and HONO production on all sites. In other words, H<sub>2</sub>O<sub>2</sub> must be a minor product, consistent with the previous experiments[97]. The  $\Delta G_{\text{RPD}}$ -limiting step for the perfect surface is OH\* dissociation(R2) and the  $\Delta G_{\text{RPD}}$ -limiting step for the defective surface is (R19)  $\text{H}_2\text{O}(\text{g}) + \text{O}(\text{Ti})^* + * \rightarrow \text{H}(\text{O})^* + \text{OOH}(\text{Ti})^*$  as shown **Fig. 5.4d, f**. For all studied sites, H<sub>2</sub>O<sub>2</sub> formation follows:

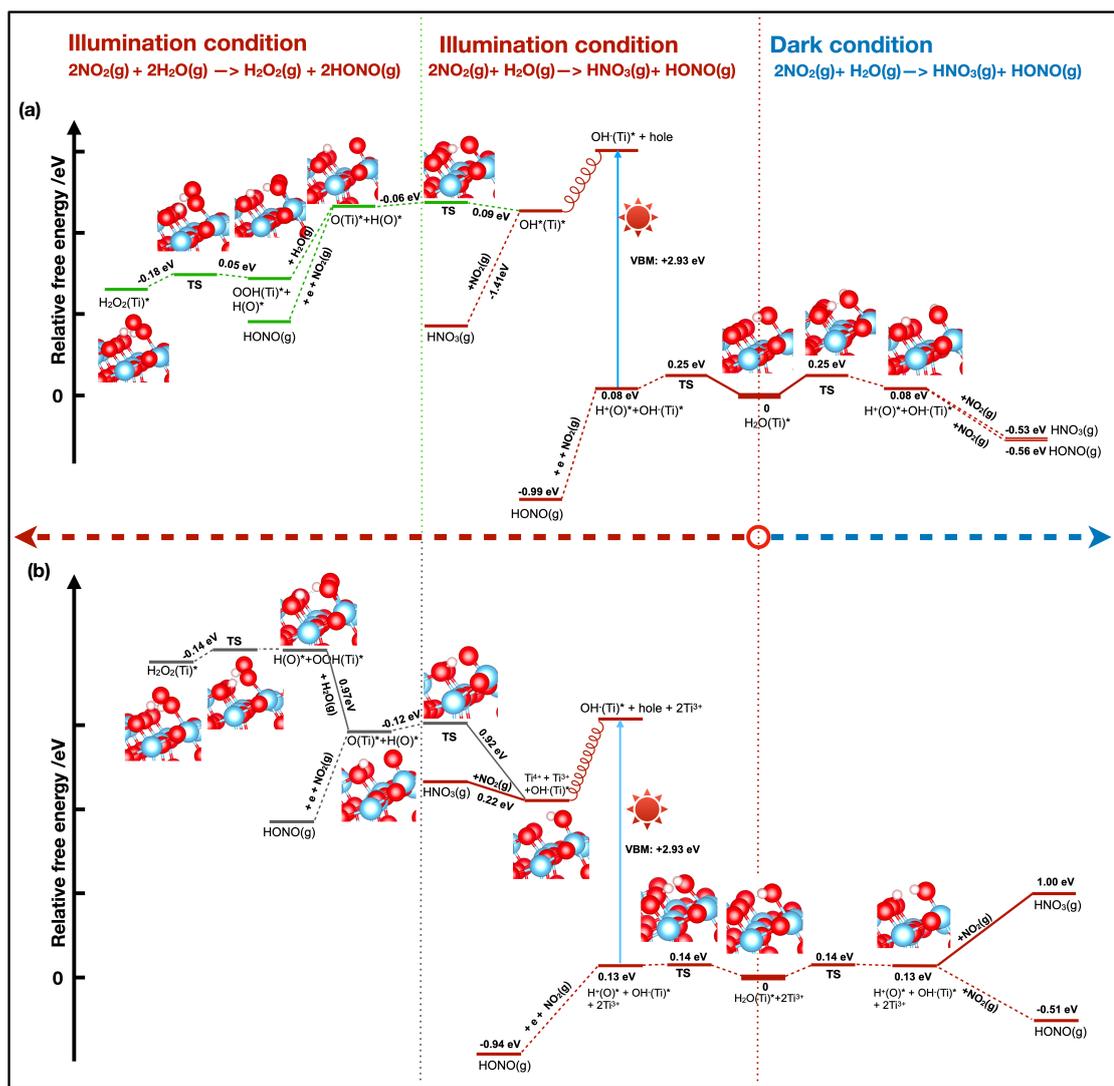


On the perfect surface, OH dissociation was facilitated by the photogenerated hole because of weakening of the bond of the adsorbed OH\* group. However, the too strong binding energy of O\* and OH\* species on the defective surface inhibited the recombinative desorption and thus reduced the activity. Note that H<sub>2</sub>O<sub>2</sub> was produced via the combination of H(O)\* and OOH(Ti)\*, instead of the previously proposed coupling of two OH\* groups[97]. Because the two Ti sites are too far apart (3.75Å) and the OH\* groups are strongly bonded.

In summary, it was found that the main products, HNO<sub>3</sub> and HONO, were identical over all sites and exhibited a high activity on the perfect surface and a lower one on the defective surface because of the strong binding of OH\* species under both dark and illumination conditions. On the defective surface, the reduced Ti (i.e., the vacancy) captured the hole rather than the OH\* group. Hence, the deNO<sub>2</sub> activities are not greatly improved on the defective surface with illumination. The main reaction pathways were not affected by the illumination. H<sub>2</sub>O<sub>2</sub> is a minor product on the perfect surface with illumination, and it was formed via the combination of OOH(Ti)\* and H(O)\*.

#### 5.4 Kinetic free energy profiles for deNO<sub>2</sub> under dark and illumination conditions

All the reaction mechanism discussions above are based on the global optimization of thermodynamics over diverse pathways. Calculated barriers are also surmountable, confirming the mechanisms above are feasible. **Fig. 5.5a** shows the reaction barriers in the production processes of HONO, HNO<sub>3</sub>, and H<sub>2</sub>O<sub>2</sub> on the perfect surface under dark and illumination conditions. Water dissociation, with a free energy barrier of 0.25 eV, and the following free energies associated with HONO and HNO<sub>3</sub> production are all downhill. In the case of H<sub>2</sub>O<sub>2</sub> formation with illumination, OH\* radical dissociation and combination of H\* and OOH\* groups are uphill with the low energy barriers of 0.09 eV and 0.05 eV, respectively. The free energies of the remaining associated steps are all downhill. The OH\* dissociation and the final steps of HONO and HNO<sub>3</sub> formation are promoted by illumination. In addition, water dissociation is a stepwise process: electron transfer (ET) follows proton transfer (PT)[117]. Namely, OH<sup>-</sup> captures the photogenerated hole after water deprotonation, yielding the OH radical. Hence, the HONO, HNO<sub>3</sub>, and H<sub>2</sub>O<sub>2</sub> production on the perfect surface is kinetically and thermodynamically favored. There are two possible reaction pathways for HONO and HNO<sub>3</sub> formation. **Fig. 5.4b** shows the optimal reaction mechanism and barriers of HONO, HNO<sub>3</sub>, and H<sub>2</sub>O<sub>2</sub> production on the defective surface, exemplified on Ti<sub>5c</sub>, under dark and illumination conditions. The dissociated water in the Eley-Rideal pathway is con-

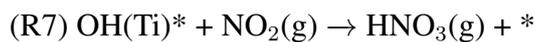


**Fig. 5.5** Calculated reaction energies and barriers of each step for the deNO<sub>2</sub> reaction mechanisms on the (a) perfect and (b) defective (Ti<sub>5c</sub>) surfaces under dark and illumination conditions, respectively. The solid lines denote the limiting steps. The VBM value was taken from experiment[8].

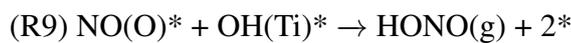
sumed by



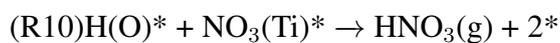
and



The second pathway includes the disproportionation of NO<sub>2</sub>, and dissociated water is consumed by

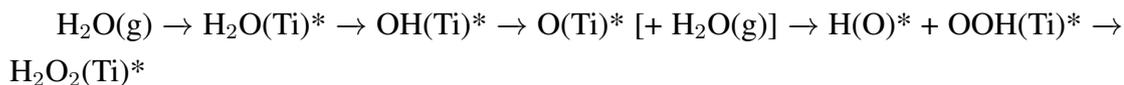


and



The energetic costs for OH(Ti)\* + NO<sub>2</sub>(g) → HNO<sub>3</sub>(g) + \* (R7) and 2NO<sub>2</sub>(g) + 2\* → NO(O)\* + NO<sub>3</sub>(Ti)\* (R8) are 0.87 eV and 1.01 eV under dark conditions, respec-

tively. In other words, disproportionation of NO<sub>2</sub> under dark conditions is a feasible pathway, which is consistent with the experiment[12]. However, HONO and HNO<sub>3</sub> production in the Eley-Rideal mechanism can be strongly facilitated by illumination, leading to the enhanced deNO<sub>2</sub> activity on Ti<sub>5c</sub>, with a barrier of 0.22 eV. Assuming NO<sub>2</sub> disproportionation is not directly affected by illumination, the two paths are both feasible with illumination. As water dissociated, the hole was captured by the reduced Ti, instead of OH\* and NO\* groups. H<sub>2</sub>O<sub>2</sub> production with illumination proceeds along the following pathway:



The dissociative barrier of OH\* groups was calculated to be 0.92 eV, which is consistent with the report published earlier[118]. The thermodynamic free energies for the O(Ti)\* + H<sub>2</sub>O(g) reaction has reached 0.97 eV. These can be ascribed to the too strong interactions of the O\* species with the defective surface, which can be explained by the fact that the excess electron is transferred to the adsorbates from the reduced Ti site. Thus, it is difficult to produce H<sub>2</sub>O<sub>2</sub>, although the combination of OOH(Ti)\* and H(Ti)\* to H<sub>2</sub>O<sub>2</sub> is a barrierless process. In short, the calculated kinetic results are consistent with the prediction from energetic global optimization of reaction phase diagram analysis.

## 5.5 Summary

In this work, I have employed an improved scheme to study the mechanisms of photoassisted deNO<sub>2</sub> reactions on the active sites of the anatase TiO<sub>2</sub>(101) surface. The electronic structures and energies were calculated by using a GGA + LZ functional, which significantly improved the accuracy (comparable to calculations with the hybrid HSE06 functional) while maintaining the efficiency of normal GGA calculations. In addition, a novel algorithm was applied to study all possible reaction pathways and to obtain a global optimization of the limiting energies. It was found that the activity on the perfect surface is higher than on a defective one, because the strong OH\* adsorption inhibits its combination with NO<sub>2</sub>. It was found the photogenerated hole is more important for the perfect surface in tuning the reactivity. Because the hole is trapped by a reduced Ti on the defective surface, it has no effect on the adsorption energy of OH\* species on the oxygen vacancy. In contrast, the OH\* adsorption energy on a perfect surface can be weakened in the presence of the photogenerated hole, resulting in enhanced OH\* conversion. This work proved the concept that the tunable reactivity in a charge-assisted process and diverse reaction pathways in photocatalysis are two methods that break the conventional scaling relation and the activity volcano plot limit under dark conditions, which is a quite critical insight for the design of catalyst.

## 6 | Activity and reaction mechanism mapping of photocatalytic deNO processes on titania surfaces

As discussed in last chapter, the photogenerated hole can promote the oxidative reactions and oxidize the  $\text{NO}_x$  into higher oxidization state with active oxygen or OH radicals[2, 7]. However, the accompanying pollution, i.e. nitrate and nitrite, is troublesome.  $\text{NO}_x$  can also be reduced with illumination[2, 4, 97]. The photogenerated electrons promote the reductive reactions and convert  $\text{NO}_x$  into harmless  $\text{N}_2$ , which can avoid secondary pollution problems[4]. Therefore, many studies have focused on the photoreduction of  $\text{NO}_x$  on  $\text{TiO}_2$  surfaces, especially the decomposition of NO into  $\text{N}_2$  and  $\text{O}_2$ , which was considered as a desirable photocatalytic denitrification process[2]. It was reported that NO can be reduced to  $\text{N}_2\text{O}$  and  $\text{N}_2$  over  $\text{TiO}_2$  powders, but  $\text{N}_2\text{O}$  is the major product[119]. Many efforts have been made to increase the selectivity of  $\text{N}_2$  in photocatalytic deNO over  $\text{TiO}_2$ [119, 120, 121, 122]. A notable finding by Wu and Krol is that oxygen vacancies are the active centers for denitrification. Increasing the number of oxygen vacancies, such as by doping the  $\text{TiO}_2$  with iron, can improve the  $\text{N}_2$  selectivity[4]. Note that oxygen vacancies are not typically present on the anatase surfaces, but they can be created by experimental techniques[4, 123]. Additionally, the reaction mechanism of photo-decomposition of NO is dispute. Wu and Krol proposed that two NO molecule can be directly dissociated to two surface oxygen and two nitrogen atoms, which can be easily coupled into  $\text{N}_2$  and  $\text{O}_2$ [4]. On the reduced rutile  $\text{TiO}_2$  (110) surface, however, Xu and Cao suggested that the trans- $(\text{NO})_2$  dimers can be formed at Vo and adjacent  $\text{Ti}_{5c}$  sites, which can then be converted to  $\text{N}_2\text{O}$ [123]. In addition to the oxygen vacancies, it is well-known crystalline structures (anatase and rutile) can be quite different in reactivity. However, no significant differences between anatase and rutile were observed in experiments. The comparison of  $\text{N}_2$  selectivity in mechanism between (anatase and rutile) titania surfaces can also provide useful guidance for catalyst design in the future.

In this work, I employed a graph theory based algorithm to study possible re-

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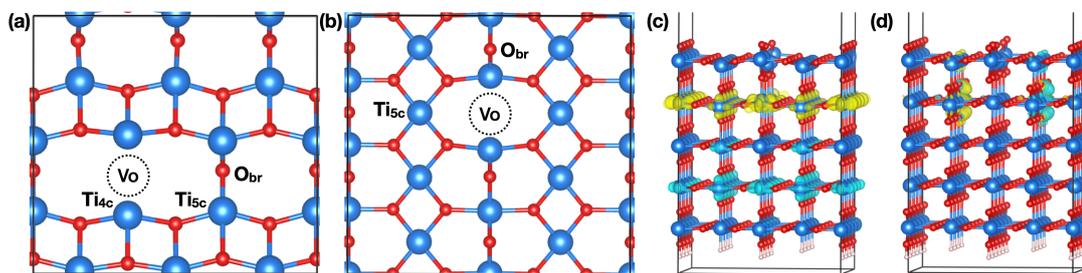
(R0) $\text{NO(g)} + * \rightarrow \text{NO(Ti)*}$
(R1) $\text{NO(Ti)*} + \text{NO(g)} \rightarrow \text{N}_2\text{O}_2\text{(Ti)*}$
(R2) $* + 2\text{NO(g)} \rightarrow \text{N}_2\text{O}_2\text{(Ti)*}$
(R3) $\text{N}_2\text{O}_2\text{(Ti)*} + * \rightarrow \text{N}_2\text{O(Ti)*} + \text{O(Ti)*}$
(R4) $2* + \text{NO(g)} \rightarrow \text{O(Ti)*} + \text{N(O)*}$
(R5) $2* + \text{NO(g)} \rightarrow \text{O(Ti)*} + \text{N(Ti)*}$
(R6) $\text{N(O)*} \rightarrow \text{N(Ti)*}$
(R7) $\text{N(O)*} + \text{N(Ti)*} \rightarrow 2* + \text{N}_2\text{(g)}$
(R8) $\text{N(Ti)*} + \text{NO(g)} \rightarrow \text{N}_2\text{O(Ti)*}$
(R9) $\text{N}_2\text{O(Ti)*} \rightarrow \text{O(Ti)*} + \text{N}_2\text{(g)}$
(R10) $\text{N}_2\text{O(Ti)*} \rightarrow \text{N}_2\text{(g)} + \text{O(O)*}$
(R11) $\text{N(O)*} + \text{NO(g)} \rightarrow \text{N}_2\text{O(O)*}$
(R12) $\text{N}_2\text{O(O)*} \rightarrow \text{N}_2\text{(g)} + \text{O(O)*}$
(R13) $\text{N}_2\text{O(O)*} \rightarrow \text{N}_2\text{O(Ti)*}$
(R14) $\text{O(O)*} \rightarrow \text{O(Ti)*}$
(R15) $2\text{N(Ti)*} \rightarrow 2* + \text{N}_2\text{(g)}$
(R16) $2\text{O(O)*} \rightarrow 2* + \text{O}_2\text{(g)}$
(R17) $2\text{N(O)*} \rightarrow 2* + \text{N}_2\text{(g)}$
(R18) $2\text{O(Ti)*} \rightarrow * + \text{O}_2\text{(Ti)*}$
(R19) $\text{O(O)*} + \text{O(Ti)*} \rightarrow * + \text{O}_2\text{(Ti)*}$
(R20) $\text{O}_2\text{(Ti)*} \rightarrow * + \text{O}_2\text{(g)}$
(R21) $\text{N}_2\text{O(Ti)*} \rightarrow * + \text{N}_2\text{O(g)}$
(R22) $\text{N}_2\text{O(O)*} \rightarrow * + \text{N}_2\text{O(g)}$
(R23) $2\text{hole} + \text{O}_2\text{(Ti)*} \rightarrow \text{Pos-O}_2\text{(Ti)*}$
(R24) $\text{Pos-O}_2\text{(Ti)*} \rightarrow * + \text{O}_2\text{(g)}$

**Table 6.1** All considered elementary steps for  $2\text{NO(g)} \rightarrow \text{N}_2\text{(g)} + \text{O}_2\text{(g)}$  and  $4\text{NO(g)} \rightarrow \text{N}_2\text{O(g)} + \text{O}_2\text{(g)}$  reactions.

action pathways[5] and obtain the collection of reaction pathways prior to energetic comparison. Then, I calculated the adsorption energies of relevant intermediates in deNO on a set of active sites over perfect and defective anatase  $\text{TiO}_2(101)$  surfaces and rutile  $\text{TiO}_2(110)$  surfaces in both dark and illumination conditions, using our recently developed density functional with appropriate band gap and localization error corrections[46], and charge correction scheme[48]. At last, the reaction phase diagrams were established based on the global (thermodynamic) energy optimization. Herein, my aim is to elucidate the deNO activity and selectivity of different active sites on the  $\text{TiO}_2$  surfaces under dark and illumination conditions.

According to the experimental observations, I considered two main reactions:  $2\text{NO(g)} \rightarrow \text{N}_2\text{(g)} + \text{O}_2\text{(g)}$  and  $4\text{NO(g)} \rightarrow \text{N}_2\text{O(g)} + \text{O}_2\text{(g)}$ . 25 possible elementary steps were considered as listed in **Table 6.1** and 3109 possible reaction pathways were obtained with the CatRPD code[124] to build the reaction networks. In short, all possible reaction pathways were obtained by the free combination of all forward and backward elementary steps. At a given active site, I first determine the most difficult steps in

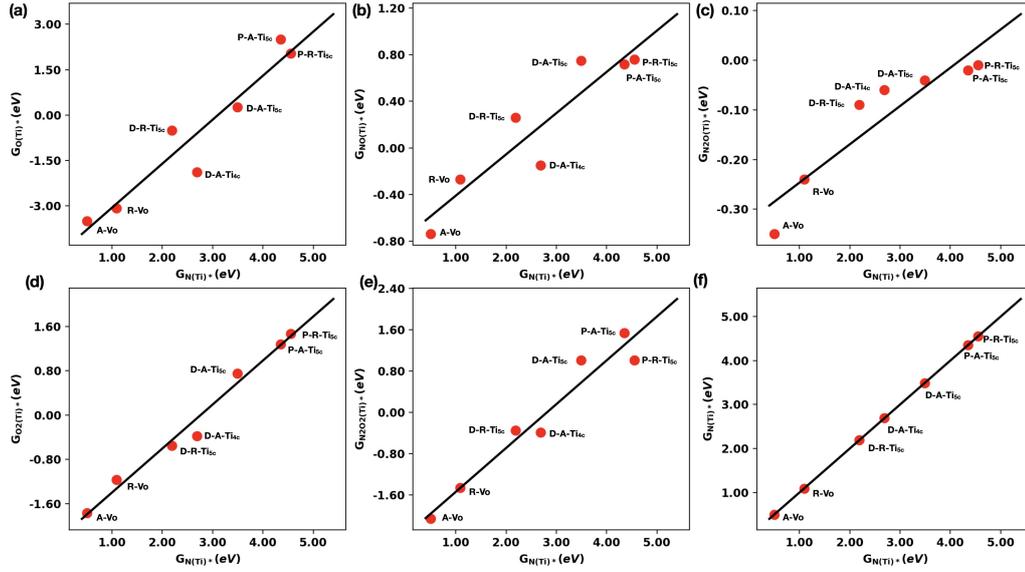
the internal comparison among elementary steps. Then, the optimal reaction pathway should be the one with minimal free energy change ( $\Delta G$ ) in external comparison among all possible pathways with Eq. 3.9.



**Fig. 6.1** Top view and nomenclature of the sites on the defective anatase  $\text{TiO}_2(101)$  surface (a) and rutile  $\text{TiO}_2(110)$  surface (b):  $\text{Vo}$  denotes the oxygen vacancy;  $\text{Ti}_{4c}$  denotes the 4-fold coordinated Ti;  $\text{Ti}_{5c}$  denotes the 5-fold coordinated surface Ti site;  $\text{O}_{br}$  denotes the bridging oxygen on the surface. Spin distributions in the singlet state of a neutral oxygen vacancy on the rutile  $\text{TiO}_2(110)$  surface, as obtained from PBE without (c) and with (d) the Lany-Zunger correction. The two spin directions are distinguished by color. Cyan, red, and pink denote Ti, O, and H, respectively.

I calculated the adsorption free energies of N, O,  $\text{N}_2\text{O}_2$ ,  $\text{N}_2\text{O}$ ,  $\text{O}_2$  on seven different active sites: oxygen vacancy site on anatase  $\text{TiO}_2(101)$  surface (A-Vo) and rutile  $\text{TiO}_2(110)$  surface (R-Vo), four-fold coordinated Ti site on the defective anatase  $\text{TiO}_2(101)$  surface (D-A- $\text{Ti}_{4c}$ ), five-fold coordinated Ti site on the defective anatase  $\text{TiO}_2(101)$  surface (D-A- $\text{Ti}_{5c}$ ) and rutile  $\text{TiO}_2(110)$  surface (D-R- $\text{Ti}_{5c}$ ), five-fold coordinated Ti site ( $\text{Ti}_{5c}$ ) on the perfect anatase  $\text{TiO}_2(101)$  surface (P-A- $\text{Ti}_{5c}$ ) and rutile  $\text{TiO}_2(110)$  surface (P-R- $\text{Ti}_{5c}$ ), respectively. (see in **Fig. 6.1a, b**) Note that I have also considered the  $\text{Ti}_{5c}(\text{Vo})$  site in my calculations but the adsorbates tend to move to the Vo site from  $\text{Ti}_{5c}(\text{Vo})$  site after the relaxation. The nomenclature is the same here for all intermediates, unless stated otherwise. As shown in **Fig. 6.1c, d**, PBE calculations show that the electrons of neutral oxygen vacancy delocalize in the whole slab while the PBE + LZ calculations show that two electrons of neutral oxygen vacancy localize on two subsurface Ti atoms, respectively. Neutral oxygen vacancy of rutile (110) surface is a biradical singlet state. These results are consistent with higher level (HSE06) calculations in the published paper[125, 76].

The adsorption free energies of N was chosen as the descriptor to establish the correlation relationships with the adsorption energies of other considered adsorbates. As shown in **Fig. 6.2**, the scaling relations show good linearity. Thus, the adsorption free energy of each considered species and the free energy of each elementary step can be well described by the descriptor.

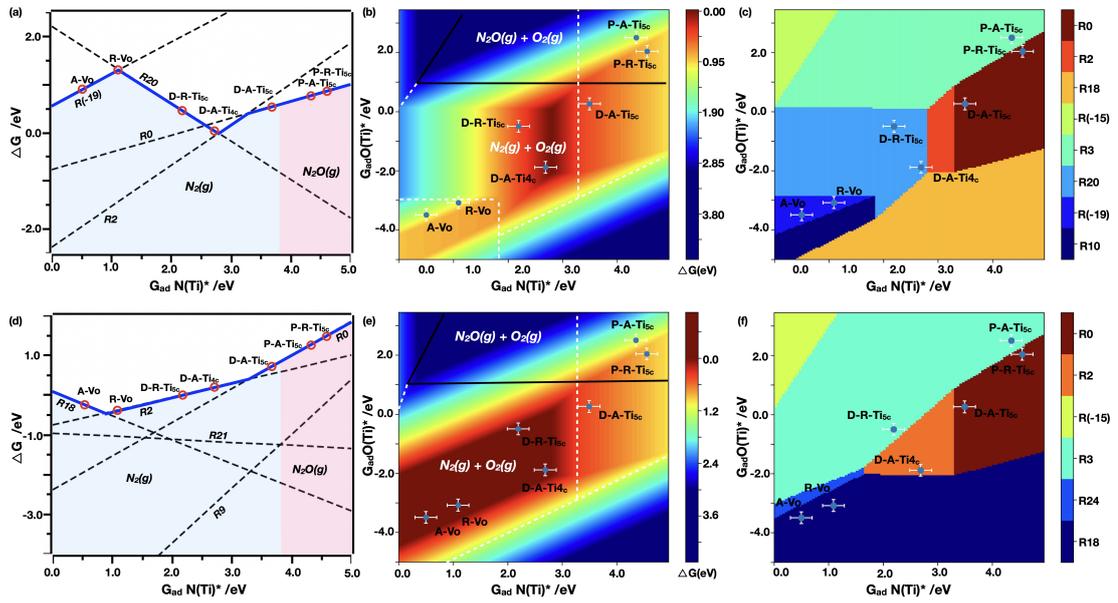


**Fig. 6.2** Scaling relations of the adsorption energies for (a)  $O^*$ , (b)  $NO^*$ , (c)  $N_2O^*$ , (d)  $O_2^*$ , (e)  $N_2O_2^*$ , (f)  $N^*$  on the different active sites of anatase  $TiO_2(101)$  and rutile  $TiO_2(110)$  surfaces. The adsorption free energies of  $N(Ti)^*$  was chosen as the descriptor. All the free adsorption energies were calculated with reference to the energy of the gas molecules  $NO$ ,  $O_2$ . A-Vo and R-Vo denote the oxygen vacancies on the surfaces of anatase  $TiO_2(101)$  and rutile  $TiO_2(110)$ , respectively. D-A- $Ti_{5c}$  and D-R- $Ti_{5c}$  denote the five-fold coordinated Ti sites on the defective surfaces of anatase  $TiO_2(101)$  and rutile  $TiO_2(110)$ , respectively. D-A- $Ti_{4c}$  denotes the four-fold coordinated Ti site on the defective anatase  $TiO_2(101)$  surface. P-A- $Ti_{5c}$  and P-R- $Ti_{5c}$  denote the five-fold coordinated Ti sites on the perfect anatase  $TiO_2(101)$  and rutile  $TiO_2(110)$  surfaces, respectively.

## 6.1 Activity and mechanism mapping of NO decomposition without illumination

As  $O^*$  and  $O_2^*$  are two key species for  $O_2$  production,  $G_{ad} O(Ti)^*$  was chosen as the second descriptor for establishing two dimensional maps. The (quasi) activity and selectivity maps of deNO reactions over diverse  $TiO_2$  surfaces/sites are shown in **Fig. 6.3**. Under dark condition, it enables to produce  $N_2$  on defective  $TiO_2$  surfaces, while it is selective for  $N_2O$  production on intact  $TiO_2$  surfaces as shown in **Fig. 6.3a, b**. D-A- $Ti_{4c}$  exhibits the highest activity to produce  $N_2$  and the main  $\Delta G_{RPD}$ -limiting steps are likely to be  $N_2O_2$  formation,  $O_2$  formation and  $O_2$  desorption as shown in **Fig. 6.3c**, because of similar  $\Delta G_{RPD}$  for these steps (around 0.01 eV). However, oxygen vacancies are inactive because a high reaction energy of around 1.20 eV. Two perfect rutile and anatase surfaces have also low activity ( $\sim 1.0$  eV uphill in reaction free energies) and the main  $\Delta G_{RPD}$ -limiting steps are  $NO$  adsorption and  $N_2O_2$  dissociation as shown in **Fig. 6.3c**. Diverse reaction pathways are compared to produce  $N_2$  and  $N_2O$  over titania surfaces. D-A- $Ti_{4c}$  and D-R- $Ti_{5c}$  follow the same reaction pathway that  $N_2$

is produced via  $N_2O^*$  dissociation:  $2NO(g) \rightarrow N_2O_2(Ti)^* \rightarrow N_2O(Ti)^* \rightarrow N_2(g) + O(Ti)^*$  (path1).  $N_2$  production on oxygen vacancies follow the similar reaction pathway:  $2NO(g) \rightarrow N_2O_2(Ti)^* \rightarrow N_2O(Ti)^* \rightarrow N_2(g) + O(O)^*$ (path2). As D-A-Ti<sub>5c</sub> is located at the boundary of path1 and path3,  $N_2$  can be produced via path1 or path3:  $NO(g) \rightarrow NO(Ti)^* \xrightarrow{+NO(g)} N_2O_2(Ti)^* \rightarrow N_2O(Ti)^* \rightarrow N_2(g)$ . At Ti<sub>5c</sub> sites on perfect surfaces,  $N_2O$  is formed via path4:  $NO(g) \rightarrow NO(Ti)^* \xrightarrow{NO(g)} N_2O_2(Ti)^* \rightarrow N_2O(Ti)^* \rightarrow N_2O(g)$ , which is consistent with the calculations by Xu et al[123]. Note that the extra electrons can migrate to adsorbates from defective surfaces[40, 57, 126, 127], especially at the oxygen vacancies. Therefore, the oxygen vacancies exhibit strong reactivity, but low catalytic activity for deNO reactions.



**Fig. 6.3** (a) The reaction free energy change of key elementary steps on all sites without illumination. The bold blue line denotes the 1D  $\Delta G_{RPD}$ -limiting steps for NO decomposition reactions on different sites, which are determined with Eq. 3.9. (b) The 2D quasi activity and selectivity map is shown with respect to the two descriptors (adsorption free energies of  $N(Ti)^*$  and  $O(Ti)^*$ ) independently without illumination. The color bar denotes the  $\Delta G_{RPD}$ -limiting energies. (c) The  $\Delta G_{RPD}$ -limiting steps for the deNO over all considered sites without illumination. (d) The 1D reaction free energy change of key elementary steps on each site with illumination. (e) The 2D quasi activity and selectivity map is shown with respect to the two descriptors (adsorption free energies of  $N(Ti)^*$  and  $O(Ti)^*$ ) independently with illumination. The color bar denotes the  $\Delta G_{RPD}$ -limiting energies. (f) The  $\Delta G_{RPD}$ -limiting steps for the deNO on all sites with illumination. A-Vo and R-Vo denote the oxygen vacancies on the anatase  $TiO_2(101)$  surface and rutile  $TiO_2(110)$ , respectively. D-A-Ti<sub>5c</sub> and D-R-Ti<sub>5c</sub> denote the five-fold coordinated Ti sites on the defective anatase  $TiO_2(101)$  and rutile  $TiO_2(110)$  surfaces, respectively. D-A-Ti<sub>4c</sub> denotes the four-fold coordinated Ti site on the defective anatase  $TiO_2(101)$  surface. P-A-Ti<sub>5c</sub> and P-R-Ti<sub>5c</sub> denote the five-fold coordinated Ti sites on the perfect anatase  $TiO_2(101)$  and rutile  $TiO_2(110)$  surfaces, respectively. The corresponding elementary steps are listed in **Table 6.1**

## 6.2 Activity and mechanism mapping of NO decomposition with illumination

In principle, there are two main possibilities to improve the activity with illumination. One is to improve the reaction free energies of limiting steps in dark condition by photogenerated carriers. The other one is to construct new reaction pathways with higher activity. As the reaction pathways were all already determined before energetic comparison, it has been a global optimum in energies in dark condition. In other words, a photo-generated carrier can affect the adsorption energies of key intermediates and break the strong correlation shown in **Fig. 6.2**. Hence, I reanalyzed the (quasi) activity and selectivity map, as well as the adsorption energies of limiting steps under illumination condition. Herein, I mainly address the activity variation of oxygen vacancies with illumination. Because the generated  $O_2^*$  species on oxygen vacancies have strong interactions with surfaces and are difficult to desorb from surface. Thus, I studied the effect of two photo-generated holes on the  $O_2$  desorption. It was found that dioxygen can be produced by the coupling of two  $O(Ti)^*$  species over most sites with illumination. It was found that the holes were trapped by the adsorbed  $O_2$  species and oxidized  $O_2^{2-}$  into  $O_2^-$  at oxygen vacancies, weakening the adsorption of  $O_2$  species.

As shown in **Fig. 6.3d, e**, the  $N_2$  selectivity is not influenced by illumination, but the activities are significantly improved by illumination on oxygen vacancies, which is consistent with experimental reports[7, 121]. It exhibits the highest activity on oxygen vacancies and the optimal reaction pathways remain unchanged on all studied sites except oxygen vacancies. On oxygen vacancies, the optimal reaction pathway changes to path1 and the main  $\Delta G_{RPD}$ -limiting step is  $O_2$  formation, which is realized via the coupling of two  $O^*$  species **Fig. 6.3f**. It should be mentioned that our method finally provides the most preferred reaction pathways based on thermodynamic analysis. It was found that path3 can produce  $N_2$  at oxygen vacancies using the same  $\Delta G_{RPD}$ -limiting step as path1. Thus,  $N_2$  can be produced at oxygen vacancies through both path1 and path3. Importantly, as shown in **Fig. 6.3d**, it was found that  $N_2O$  desorption (R21) is comparable with  $N_2O$  dissociation (R9). Thus, it is possible to produce  $N_2O$  on defective surfaces with the same  $\Delta G_{RPD}$ -limiting steps as  $N_2$  production. In other words, both  $N_2$  and  $N_2O$  can be produced with similar activity on defective surfaces. In addition, it was found that  $N_2$  and  $N_2O$  are produced via the dissociation of  $N_2O_2^*$ , instead of the direct NO dissociation. As the NO dissociative energy is too high,  $\sim 1.16$  eV, on A-Vo site ( $NO(Vo)^* + * \rightarrow N(Vo)^* + O(Ti)^*$ ).

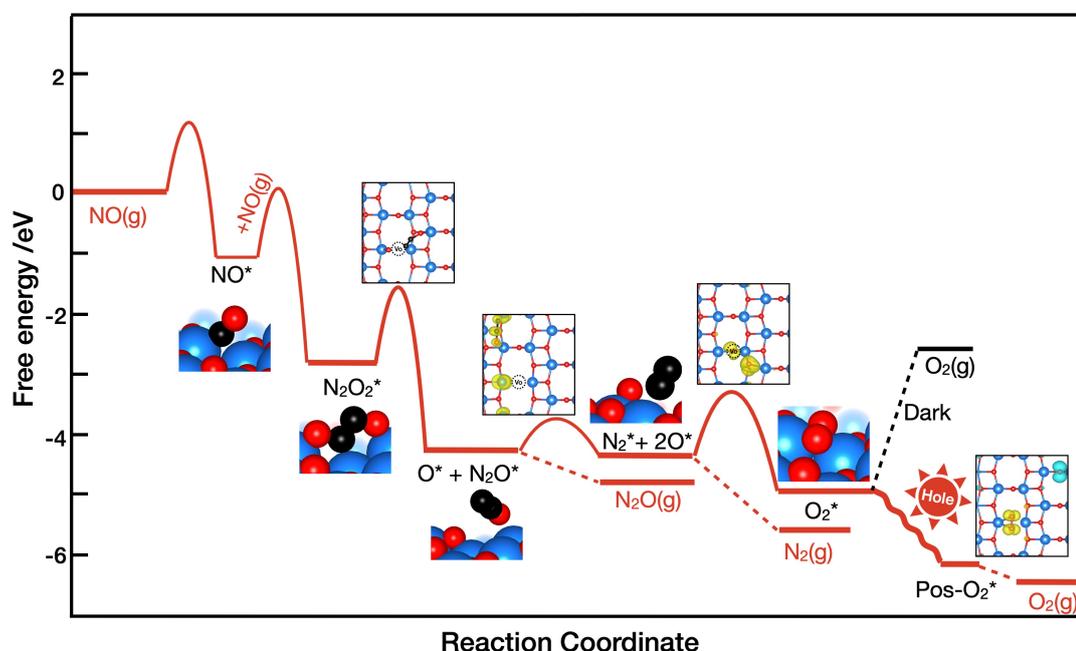
### 6.3 Summary

In a short summary, it is thermodynamically feasible to produce  $\text{N}_2$  on defective  $\text{TiO}_2$  surfaces, whereas it is more selective to produce  $\text{N}_2\text{O}$  on perfect  $\text{TiO}_2$  surfaces. The (quasi) activity was significantly enhanced by photo-generated holes and the photocatalytic decomposition of NO into  $\text{N}_2$  is thermodynamically feasible on defective titania surfaces. The low efficiency observed in experiments could be attributed to the weak adsorption of  $\text{N}_2\text{O}$ , which make it easy to desorb from surfaces rather than  $\text{N}_2\text{O}^*$  dissociation to  $\text{N}_2$ . Notably,  $\text{N}_2$  production is favored by  $\text{N}_2\text{O}$  dissociation, rather than the paths of direct NO dissociation. This is an important finding that  $\text{N}_2$  production has to via  $\text{N}_2\text{O}^*$  path. The thermodynamic analysis based on reaction phase diagram and global energy optimization provides a fundamental limit of enhancing  $\text{N}_2$  selectivity. In the following, I performed kinetic studies for photocatalytic NO decomposition over titania surfaces.

# 7 | A fundamental limit of selectivity in photocatalytic denitrification over titania surfaces

## 7.1 Microkinetic modeling of photocatalytic deNO over titania surfaces

Based on the above discussions, I mainly focused on the kinetics of photocatalytic deNO on Vo site of anatase TiO<sub>2</sub>(101) surface, because oxygen vacancy was usually found with relatively high thermodynamic feasibility for N<sub>2</sub> selectivity. A systematic microkinetic model was conducted based on DFT calculations (as shown in **Fig. 7.1**) and experimental conditions to study the rate-limiting steps and selectivity of products. The detailed input parameters for microkinetic modelling are listed in **Table 7.1**.



**Fig. 7.1** Free energy profile for photocatalytic decomposition of NO on the oxygen vacancy of anatase TiO<sub>2</sub>(101) surface. The spin densities of the key transition states are illustrated. Blue, red and black denote Ti, O and N atoms, respectively. Pos-O<sub>2</sub>\* denotes the O<sub>2</sub> adsorption in positively charged system.

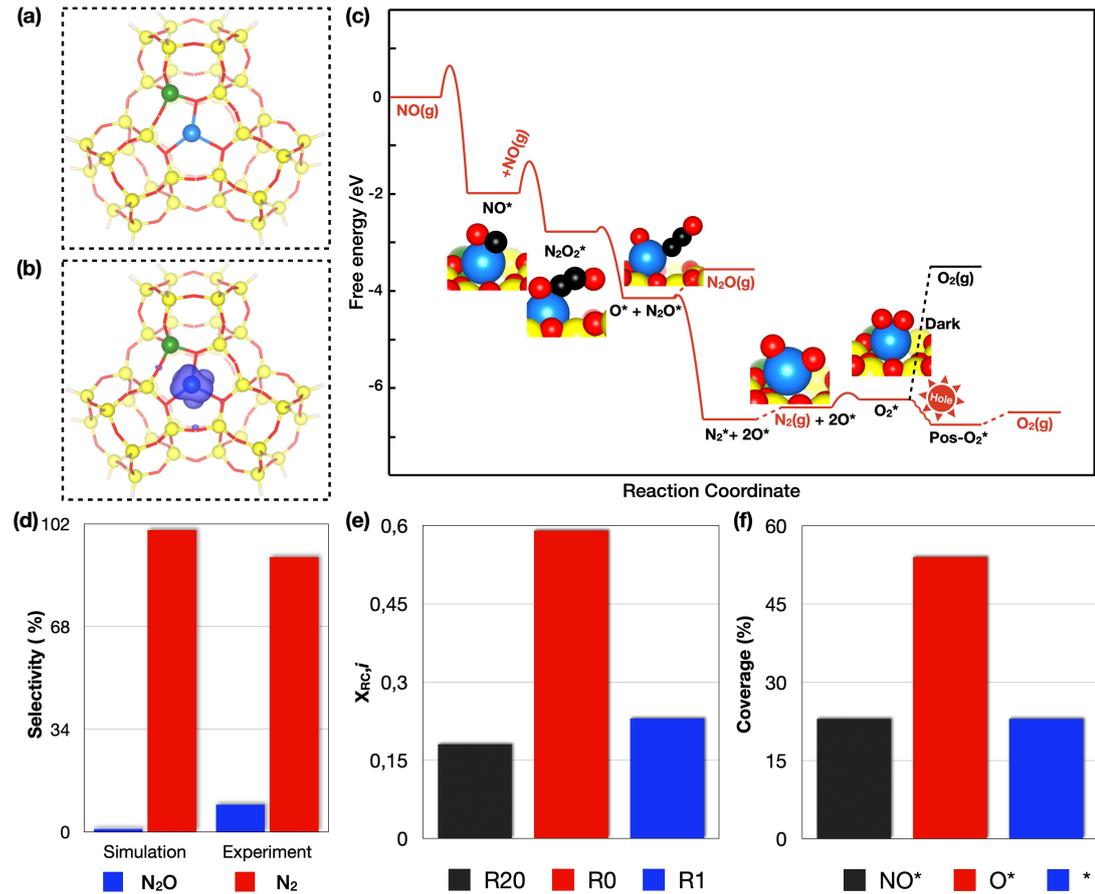
elementary steps	Zeolite(298K)		A-Vo(493K)	
	Ea / eV	$\Delta H$ / eV	Ea / eV	$\Delta H$ / eV
$\text{NO}(\text{g}) + * \rightarrow \text{NO}^*$		-2.63		-2.24
$\text{NO}^* + \text{NO}(\text{g}) \rightarrow \text{N}_2\text{O}_2^*$		-1.45		-2.91
$2\text{NO}(\text{g}) + 2* \rightarrow \text{N}_2\text{O}_2^*$		-4.08		-5.15
$\text{N}_2\text{O}_2^* + * \rightarrow \text{N}_2\text{O}^* + \text{O}^*$	0.09	-1.37		
$\text{N}_2\text{O}^* + * \rightarrow \text{N}_2^* + \text{O}^*$	0.04	-2.50		
$2\text{O}^* \rightarrow \text{O}_2^*$	0.30	0.17		
$\text{N}_2\text{O}_2^* + \# \rightarrow \text{N}_2\text{O}\# + \text{O}^*$			1.25	-1.44
$\text{N}_2\text{O}\# + \# \rightarrow \text{N}_2\# + \text{O}\#$			0.58	-0.09
$\text{O}^* + \text{O}\# \rightarrow \text{O}_2^* + \#$			1.04	-0.59
$\text{O}_2^* \rightarrow \text{O}_2(\text{g}) + *$		2.76		2.68
$\text{O}_2^* + 2\text{hole} \rightarrow \text{Pos-O}_2^*$	-	-	-	-
$\text{Pos-O}_2^* \rightarrow \text{O}_2(\text{g}) + *$		0.87		0.83
$\text{N}_2^* \rightarrow \text{N}_2(\text{g}) + *$		0.84		
$\text{N}_2\# \rightarrow \text{N}_2(\text{g}) + \#$				0.43
$\text{N}_2\text{O}^* \rightarrow \text{N}_2\text{O}(\text{g}) + *$		1.28		
$\text{N}_2\text{O}\# \rightarrow \text{N}_2\text{O}(\text{g}) + \#$				0.68

**Table 7.1** Thermodynamic and kinetic data of deNO reactions on the Ti site of the modified zeolite and Vo site on the anatase  $\text{TiO}_2(101)$  surface. # and \* denote the different active sites.

As shown in **Fig. 7.1**, a trans- $\text{N}_2\text{O}_2^*$  could be formed via the coupling of two gas phase NO molecules with energy barrier of 2.31 eV (path1) or via Eley-Rideal (E-R) mechanism with two lower energy barriers(1.15 eV) (path3). Then,  $\text{N}_2$  is formed in such a way:  $\text{N}_2\text{O}_2(\text{Ti})^* \rightarrow \text{N}_2\text{O}(\text{Ti})^* \rightarrow \text{N}_2(\text{g})$ .  $\text{O}_2^*$  is formed via the coupling of two  $\text{O}^*$  species and the formed  $\text{O}_2^*$  has strong interaction with the catalyst. Under illumination condition, the photo-generated holes weaken the  $\text{O}_2^*$  adsorption and the photo-generated electrons are supposed to reduce the oxygen vacancy. The optimal reaction path was indicated by the red lines in **Fig. 7.1**. The free energy barriers associated with  $\text{N}_2$  production are uphill. It is found that the main rate determining steps are  $\text{N}_2\text{O}_2^*$  formation,  $\text{N}_2\text{O}_2^*$  dissociation and  $\text{O}_2^*$  formation with energy barriers of 1.15 eV, 1.25 eV and 1.04 eV, respectively. Notably,  $\text{N}_2\text{O}$  is favored to desorb in comparison with  $\text{N}_2\text{O}$  dissociation. Therefore, the competing  $\text{N}_2\text{O}$  production and the high kinetic barriers of rate limiting steps lead to the inefficient  $\text{N}_2$  production at A-Vo site. This is generally consistent with the experimental results that the major product is  $\text{N}_2\text{O}$ [119]. I have analyzed the electronic structures in photocatalytic processes and the spin densities of key transition states are shown in **Fig. 7.1**. The adsorbed  $\text{N}_2\text{O}_2^*$  has a binding with Vo and  $\text{Ti}_{5c}$  and captures the electrons from oxygen vacancy, weakening the N-O bond. In transition state of  $\text{N}_2\text{O}$  dissociation, an electron transfers to  $\text{N}_2\text{O}$  from oxygen vacancy, causing bending of liner  $\text{N}_2\text{O}$  and shortening the Ti-O and Ti-N bonds from 2.80 Å to 2.35 Å, facilitating N-O bond breaking. The two extra electrons in oxygen

vacancy are trapped by two O species in the transition state of  $O_2$  formation. Finally, the formed  $O_2^{2-}$  is oxide to  $O_2^-$  by the photogenerated hole. Therefore, the oxygen vacancy facilitates the bond breaking and increases the reactivity under illumination condition.

## 7.2 Microkinetic modeling of photocatalytic deNO on Ti modified Y-zeolite



**Fig. 7.2** (a) Ti modified Y-zeolite cluster model. (b) Spin density distributions in the Ti-modified Y-zeolite cluster with isovalue of  $0.005 a_0^{-3}$  ( $a_0$ : Bohr radius). (c) Energy profile for photo-decomposition of NO on Ti site of modified zeolite. (d) The simulated and experimental selectivities in photocatalytic decomposition of NO on Ti-modified zeolite. (e) The degree of rate control (DRC) analysis in the photocatalytic deNO processes for  $N_2$  production on zeolite. (f) The coverage of species distribution at steady state. Pos- $O_2^*$  denotes the  $O_2$  adsorption in positively charged system. \* denotes the active site Ti of zeolite. Cyan, red, yellow, green and pink denote Ti, O, Si, Al and H atoms, respectively. The corresponding elementary steps are listed **Table 6.1**.

It was observed that the maximum conversion of NO is only 1% on the reduced  $TiO_2$  and 4.5% on the Fe-doped  $TiO_2$ , respectively[4]. The observed activity and selectivity is still low, so it is crucial to find the chemical origins of the low efficiency of NO photo-decomposition. In contrast, there is an exceptional success found in Ti-modified zeolite,

where the N<sub>2</sub> selectivity is much improved (>90%)[121]. It can be quite helpful in the choice of experiments as far as the fundamental limit of photocatalytic denitrification on titania surfaces was understood.

elementary step(i)	r <sub>i</sub>	Z <sub>i</sub>	X <sub>RC,i</sub>
NO(g) + * → NO*	1.38e-05	3.39e-28	0.59
NO* + NO(g) → N <sub>2</sub> O <sub>2</sub> *	1.38e-05	4.42e-23	0.23
N <sub>2</sub> O <sub>2</sub> * + * → N <sub>2</sub> O* + O*	1.38e-05	2.33e-24	0.00
N <sub>2</sub> O* + * → N <sub>2</sub> * + O*	1.38e-05	3.58e-29	0.00
2O* → O <sub>2</sub> *	1.38e-05	1.00e-01	0.00
O <sub>2</sub> * + 2hole → Pos-O <sub>2</sub> *	-	-	-
Pos-O <sub>2</sub> * → O <sub>2</sub> (g) + *	1.38e-05	3.47e-01	0.18
N <sub>2</sub> * → N <sub>2</sub> (g) + *	1.38e-05	7.30e-01	0.00

**Table 7.2** Kinetic Data, r<sub>i</sub>, the reversibility (Z<sub>i</sub>), and X<sub>RC,i</sub>, for each elementary reaction step in the deNO over Ti modified zeolite. \* denotes the Ti site of zeolite. Note that the rate r<sub>i</sub> is equal to the difference between the forward (r<sub>i</sub><sup>+</sup>) and reverse (r<sub>i</sub><sup>-</sup>) rates, namely, r<sub>i</sub> = r<sub>i</sub><sup>+</sup> - r<sub>i</sub><sup>-</sup>

The unique performance of Ti-modified zeolite frameworks, where the photocatalysts were prepared by ion exchange techniques[121], was ascribed to the observed highly dispersed molecular-sized titanium oxide species incorporated within zeolite frameworks, which was not observed in TiO<sub>2</sub> powder by ESR and XAFS measurements [120, 121]. To figure out the distinctive behaviors on Ti-modified zeolites, I built a Y-zeolite model, where a Si atom is replaced by a Al atom and a Ti atom is introduced in the center of six-ring. Al atom takes one electron from Ti atom and oxides Ti<sup>0</sup> into Ti<sup>+</sup>. The side view and spin density of Ti modified Y-zeolite are shown in **Fig. 7.2a, b**. A microkinetic model at steady state was performed to study the key factors of N<sub>2</sub> selectivity, compared to N<sub>2</sub>O production. A comparative analysis was made to understand the exceptional success of improved N<sub>2</sub> selectivity in Ti-modified zeolite with respect to titania surfaces. It was found the adsorption energies of N<sub>2</sub>O must be strengthened to enhance the N<sub>2</sub> selectivity. The details for microkinetic modelling are listed in **Table 7.1**. The calculated free energy diagram is shown in **Fig. 7.2c**. The reaction pathway is consistent with that of the A-Vo. The free energies associated with N<sub>2</sub> production are barrierless except the N<sub>2</sub>O<sub>2</sub>\* formation. N<sub>2</sub>O<sub>2</sub>\* prefers to form via E-R mechanism with two lower energy barriers of 0.65 eV. Note that N<sub>2</sub>O\* adsorption is stronger than that on titania surfaces and N<sub>2</sub>O desorption is uphill on Ti-modified Y-zeolite. The intrinsic activity and selectivity of N<sub>2</sub> and N<sub>2</sub>O, and the coverage of intermediates at steady state are summarized in **Fig. 7.2d, e**. The simulated N<sub>2</sub> selectivity on the Ti-modified zeolite is up to 100%, as shown in **Fig. 7.2d**, which is basically consistent with the experimental report of 91%[121]. The DRC (degree of rate control) X<sub>RC,i</sub> (see the definition in

ref.[93]) was usually used to analyze the importance of each elementary step  $i$  in the studied mechanism[20, 92, 93]. The larger  $X_{RC,i}$  value of a given elementary step  $i$ , the greater effect it has on the whole reaction rate. Namely, the elementary step with largest  $X_{RC,i}$  in the proposed mechanism is the rate determining step. The DRC of elementary steps R0 (NO adsorption), R1 (the association of  $\text{NO}(\text{Ti})^*$  and  $\text{NO}(\text{g})$ ) and R20 ( $\text{O}_2$  desorption) are 0.59, 0.23, 0.18 in the proposed photocatalytic deNO reaction mechanisms, respectively (**Fig. 7.2e and Table 7.2**). The DRC are negligibly small for all the other elementary steps. Thus, the R0 (NO adsorption) is the main kinetically rate limiting step in photocatalytic deNO processes on Ti-modified Y-zeolite. As shown in **Fig. 7.2f**, the coverage of  $\text{O}^*$ , free active site, and  $\text{NO}^*$  are 54%, 23%, 23%, respectively. This implies  $\text{O}^*$  species have a relatively strong adsorption and they are hard to be coupled to dioxygen. The desorption of  $\text{NO}^*$  species is suppressed and there are also some available free active sites. Therefore, the  $\text{N}_2$  selectivity and yield can be further increased by weakening  $\text{O}^*$  adsorption and increasing  $\text{N}_2\text{O}^*$  adsorption and dissociation. This should be a new strategy for photocatalyst design in the future.

### 7.3 Summary

Microkinetic modeling was applied to understand the kinetic factors at experimental conditions, clarifying the fundamental limits in deNO on titania surfaces. Although it was found that the deNO activity can be significantly enhanced under illumination condition by weakening the  $\text{O}_2$  adsorption, the limiting factor is that the  $\text{N}_2\text{O}$  and  $\text{N}_2$  productions share the same  $\text{N}_2\text{O}^*$  species. A exceptional character of Ti-modified zeolites with good  $\text{N}_2$  selectivity can be attributed to very strong  $\text{N}_2\text{O}^*$  adsorption, compared to the weak  $\text{N}_2\text{O}$  adsorption and low  $\text{N}_2$  selectivity over titania surfaces. Therefore,  $\text{N}_2\text{O}^*$  is a key intermediate to enhance  $\text{N}_2$  selectivity and deNO activity. I demonstrate here the reaction phase diagram analysis can be a good tool to obtain a global picture of activity (and selectivity) trend over different reaction sites. Combined with microkinetic modelling, the key factors of catalytic failure and exceptional success can be rationalized well with relatively low computational cost, which can be a good choice for catalyst design.

## 8 | Conclusion

In this work, I have adopted an improved scheme to study the mechanisms of photo assisted deNO<sub>x</sub> reactions on the active sites of the various titania surfaces, perfect and defective anatase TiO<sub>2</sub>(101) and rutile TiO<sub>2</sub>(110) surfaces. On the one hand, the electronic structures and energies were calculated using the GGA + LZ functional, which greatly improved the accuracy (comparable to the calculations using the hybrid HSE06 functional) while maintaining the efficiency of normal GGA calculations. On the other hand, a novel algorithm was applied to study all possible reaction pathways and to obtain a global optimization of the limiting energies. Additionally, microkinetic modeling was applied to understand the kinetic factors under experimental conditions, clarifying the fundamental limits in deNO over titania surfaces.

1. I investigated the reaction mechanisms of NO<sub>2</sub> and H<sub>2</sub>O on anatase TiO<sub>2</sub>(101) surface. It was found that the major products are HONO and HNO<sub>3</sub> and the minor product is H<sub>2</sub>O<sub>2</sub>. The activity is higher on the perfect surface than on a defective surface. Because the strong OH\* adsorption inhibits its association with NO<sub>2</sub> on the defective surface. The photogenerated carriers favor the tuning of the reactivity of perfect surface while having no effect on the reactivity of defective surface. Because the OH\* adsorption energy on the perfect surface can be weakened in the presence of the photogenerated hole, resulting in enhanced OH\* conversion. In contrast, the hole is trapped by a reduced Ti on the defective surface and thus the adsorption energy of OH\* species on the oxygen vacancy remains unchanged. (Published in JPCL)
2. I have systematically investigated the thermodynamics of photocatalytic decomposition of NO (deNO) into N<sub>2</sub>O, N<sub>2</sub> and O<sub>2</sub> on a variety of active sites of anatase and rutile TiO<sub>2</sub> surfaces. The photogenerated carries can significantly improve the activity. It was found that it is selective to N<sub>2</sub> production on defective TiO<sub>2</sub> surfaces, whereas it is more selective to N<sub>2</sub>O production on perfect TiO<sub>2</sub> surfaces. N<sub>2</sub> production shared the same N<sub>2</sub>O\* species with N<sub>2</sub>O production and N<sub>2</sub> was not produced via NO direct dissociation. (Under review in JPCL)
3. Ti-modified zeolite performed distinctive N<sub>2</sub> selectivity in comparison to titania

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surfaces. Microkinetic simulations of Ti-modified zeolite and TiO<sub>2</sub> revealed that the strong adsorption of N<sub>2</sub>O\* led to the good selectivity for N<sub>2</sub> production. Consequently, the strong adsorption of N<sub>2</sub>O plays an important role in improving N<sub>2</sub> selectivity and denitrification activity. (Under review in JPCL)

It was indicated that the reaction phase diagram analysis can be a useful tool to obtain a global picture of activity (and selectivity) trend over various active sites, while the tunable reactivity in a charge-assisted process can be achieved by breaking the conventional scaling relationships and the limits of activity volcano plot under dark conditions. In combination with microkinetic modelling, the key factors for catalytic failure and exceptional success can be well rationalized at relatively low computational cost, which can be an important option for catalyst design.

## LIST OF PUBLICATIONS

1. Guo, P.; Fu, X.; Deák, P.; Frauenheim, T.; Xiao, J. Activity and Mechanism Mapping of Photocatalytic NO<sub>2</sub> Conversion on the Anatase TiO<sub>2</sub>(101) Surface. *J. Phys. Chem. Lett.* 2021, 12 (32), 7708–7716.
2. Guo, P.; Deák, P.; Fu, X.; Frauenheim, T.; Xiao, J. A fundamental limit of selectivity in photocatalytic denitrification over titania, *J. Phys. Chem. Lett.* 2022, 13, 11051–11058.

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# Appendices

# A | Total energies corrections for charged systems

System with a positive charge	Adsorption energies (corrections)	
	SLABCC	SCPC
$\text{Ti}_{5c}^{perfect} + \text{NO}$	2.48 eV (-0.01 eV)	2.54 eV (0.05 eV)
$\text{O}_{br}^{perfect} + \text{NO}$	1.56 eV (-0.03 eV)	1.62 eV (0.03 eV)
$\text{Ti}_{5c}^{defective} + \text{NO}$	2.52 eV (0.00 eV)	2.70 eV (0.18 eV)
$\text{O}_{br}^{defective} + \text{NO}$	2.89 eV (0.00 eV)	2.96 eV (0.07 eV)
$\text{Ti}_{5c}^{perfect} + \text{OH}$	0.37 eV (-0.03 eV)	0.37 eV (-0.03 eV)
$\text{Ti}_{5c}^{defective} + \text{OH}$	-1.32 eV (-0.11 eV)	-1.31 eV (-0.10 eV)

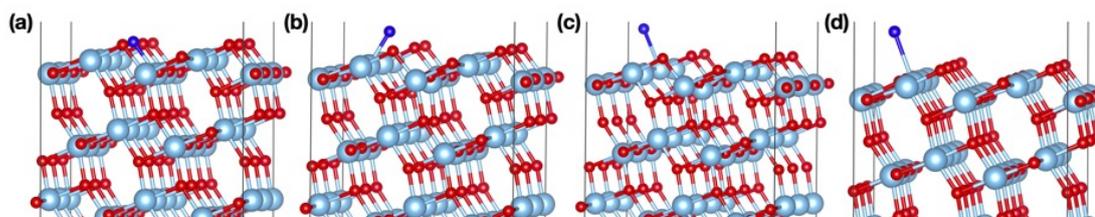
**Table A.1** Charge correction to adsorption energies with SLABCC and SCPC methods in deNO<sub>2</sub> process.

Active sites	Free adsorption energy (eV)	
	without corrections	with SCPC correction
A-Ti <sub>5c</sub> <sup>perfect</sup>	1.28	1.48 (0.20)
R-Ti <sub>5c</sub> <sup>perfect</sup>	1.21	1.04 (-0.17)
A-Vo	1.06	0.86 (-0.20)
A-Ti <sub>4c</sub> <sup>defective</sup>	1.44	0.98 (-0.46)
A-Ti <sub>5c</sub> <sup>defective</sup>	1.48	1.74 (0.26)
R-Vo	1.35	0.91 (-0.44)
R-Ti <sub>5c</sub> <sup>defective</sup>	1.53	1.65 (0.12)

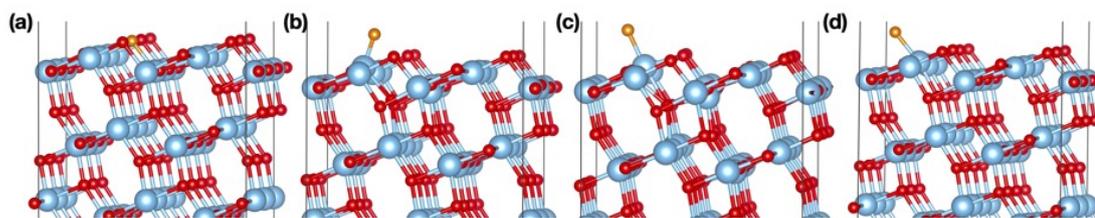
**Table A.2** The corrected adsorption energies of O<sub>2</sub> (without ZPE and entropy corrections) over different sites of TiO<sub>2</sub> surfaces in deNO process.

## B | Optimized geometries in the deNO<sub>2</sub> process

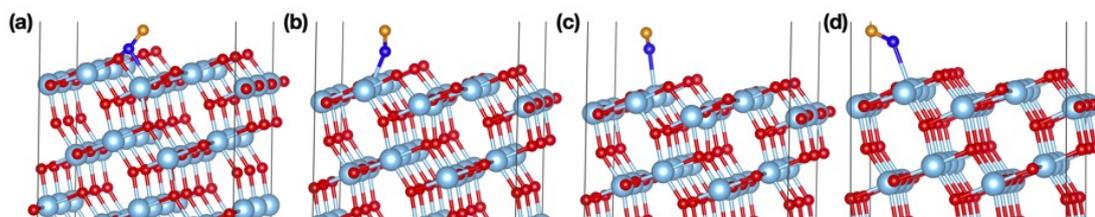
Here we display the optimized geometries of the N\*, NO\*, NO<sub>2</sub>\*, NO<sub>3</sub>\*, O\*, O<sub>2</sub>\*, OH\*, OOH\*, H<sub>2</sub>O\*, H<sub>2</sub>O<sub>2</sub>\* species adsorbed on the Vo, Ti<sub>4c</sub>, Ti<sub>5c</sub> sites of the defective anatase TiO<sub>2</sub>(101) surface and on the Ti<sub>5c</sub> site of the perfect anatase TiO<sub>2</sub>(101) surface. We have also considered the Ti<sub>5c</sub>(Vo) site in our calculations but the adsorbates tend to move to the Vo site from Ti<sub>5c</sub>(Vo) site during the relaxation.



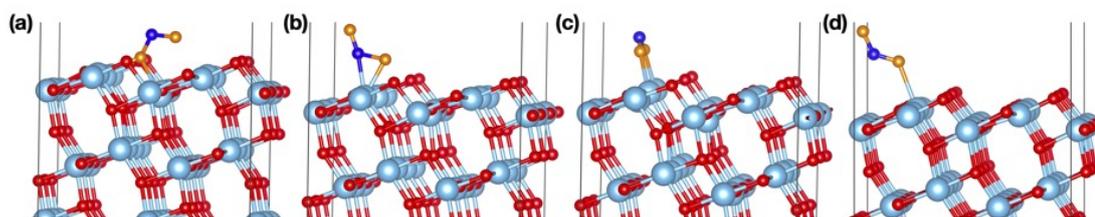
**Fig. B.1** The optimized geometries of N adsorbed on the (a) Vo, (b) Ti<sub>4c</sub>, (c) Ti<sub>5c</sub> sites of the defective surface and (d) on Ti<sub>5c</sub> of the perfect surface. Cyan, red, blue denotes Ti, O and N, respectively.



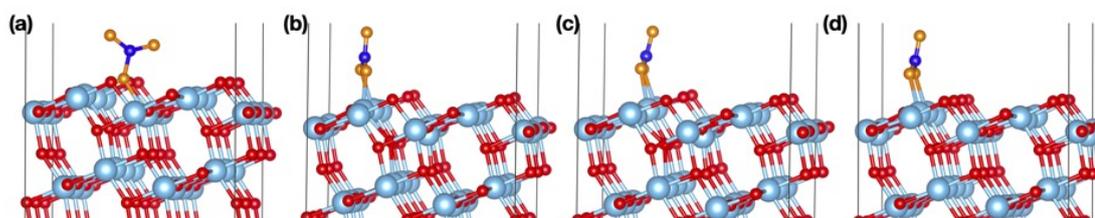
**Fig. B.2** The optimized geometries of O adsorbed on the (a) Vo, (b) Ti<sub>4c</sub>, (c) Ti<sub>5c</sub> sites of the defective surface and (d) on Ti<sub>5c</sub> of the perfect surface. Cyan, red, blue denotes Ti, O and N, respectively.



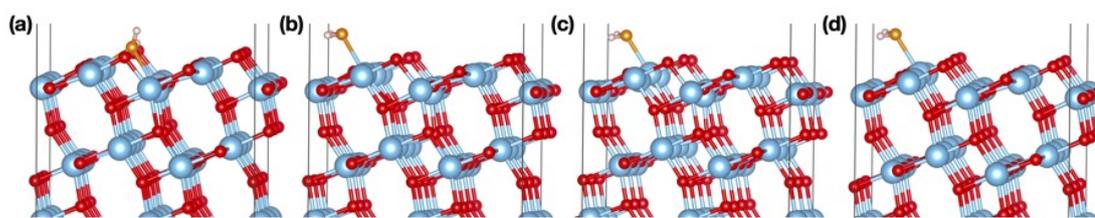
**Fig. B.3** The optimized geometries of NO adsorbed on the (a) Vo, (b) Ti<sub>4c</sub>, (c) Ti<sub>5c</sub> sites of the defective surface and (d) on Ti<sub>5c</sub> of the perfect surface. Cyan, red, blue denotes Ti, O and N, respectively.



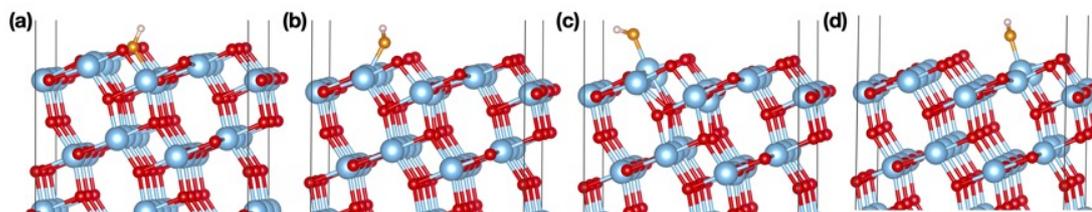
**Fig. B.4** The optimized geometries of NO<sub>2</sub> adsorbed on the (a) Vo, (b) Ti<sub>4c</sub>, (c) Ti<sub>5c</sub> sites of the defective surface and (d) on Ti<sub>5c</sub> of the perfect surface. Cyan, red, blue denotes Ti, O and N, respectively.



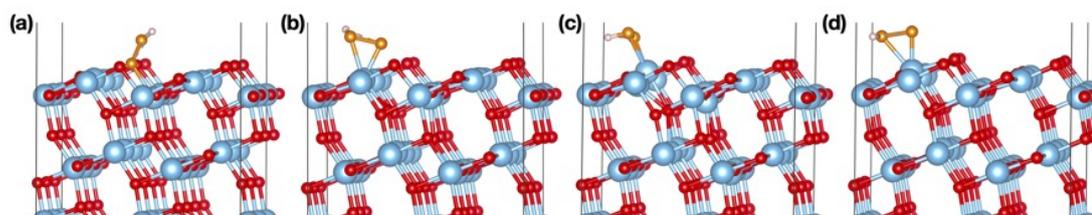
**Fig. B.5** The optimized geometries of NO<sub>3</sub> adsorbed on the (a) Vo, (b) Ti<sub>4c</sub>, (c) Ti<sub>5c</sub> sites of the defective surface and (d) on Ti<sub>5c</sub> of the perfect surface. Cyan, red, blue denotes Ti, O and N, respectively.



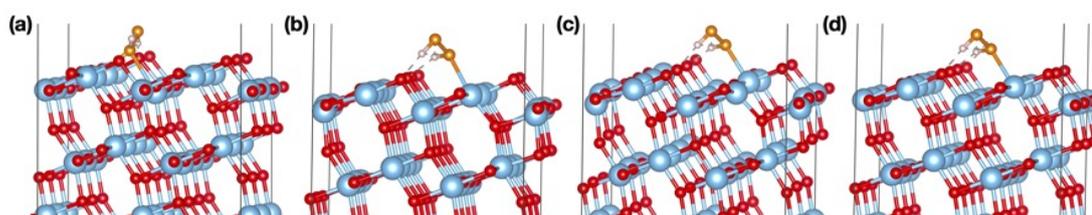
**Fig. B.6** The optimized geometries of H<sub>2</sub>O adsorbed on the (a) Vo, (b) Ti<sub>4c</sub>, (c) Ti<sub>5c</sub> sites of the defective surface and (d) on Ti<sub>5c</sub> of the perfect surface. Cyan, red, blue denotes Ti, O and N, respectively.



**Fig. B.7** The optimized geometries of OH adsorbed on the (a)  $\text{Vo}$ , (b)  $\text{Ti}_{4c}$ , (c)  $\text{Ti}_{5c}$  sites of the defective surface and (d) on  $\text{Ti}_{5c}$  of the perfect surface. Cyan, red, blue denotes Ti, O and N, respectively.



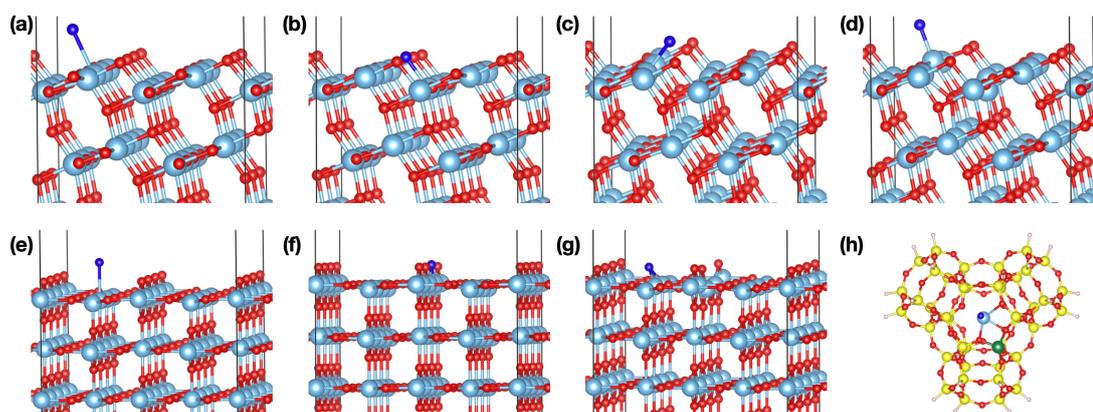
**Fig. B.8** The optimized geometries of OOH adsorbed on the (a)  $\text{Vo}$ , (b)  $\text{Ti}_{4c}$ , (c)  $\text{Ti}_{5c}$  sites of the defective surface and (d) on  $\text{Ti}_{5c}$  of the perfect surface. Cyan, red, blue denotes Ti, O and N, respectively.



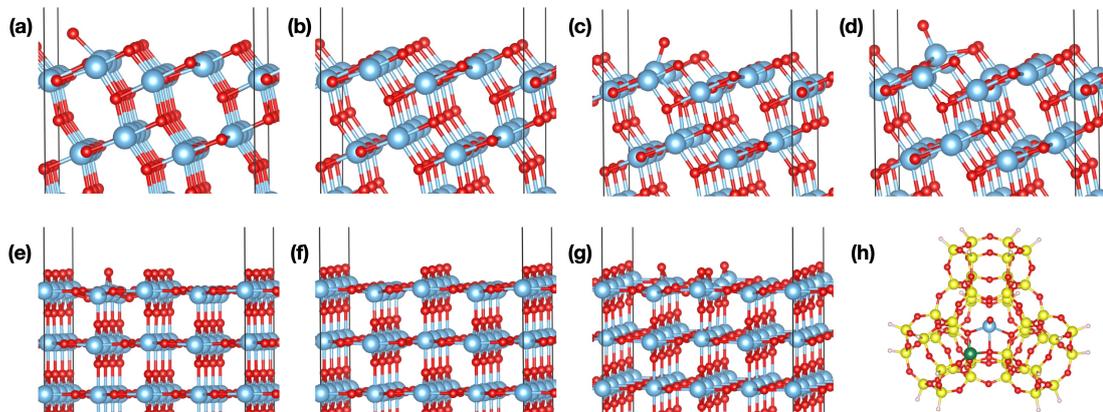
**Fig. B.9** The optimized geometries of  $\text{H}_2\text{O}_2$  adsorbed on the (a)  $\text{Vo}$ , (b)  $\text{Ti}_{4c}$ , (c)  $\text{Ti}_{5c}$  sites of the defective surface and (d) on  $\text{Ti}_{5c}$  of the perfect surface. Cyan, red, blue denotes Ti, O and N, respectively.

## C | Optimized geometries in deNO process

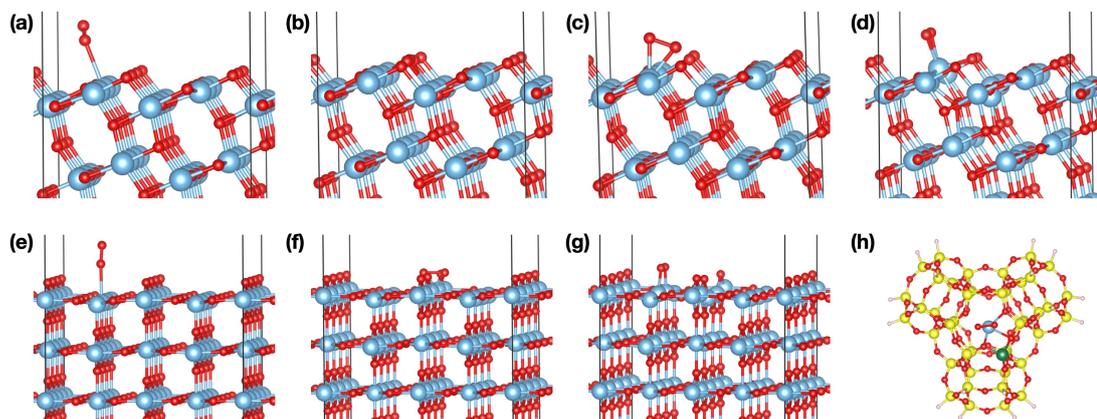
Here we displayed the optimized geometries of the  $N^*$ ,  $O^*$ ,  $O_2^*$ ,  $N_2O^*$ ,  $N_2O_2^*$  species adsorbed on the different sites of perfect and defective  $TiO_2$  surfaces and Ti site of Y-zeolite(Ti-oxide).



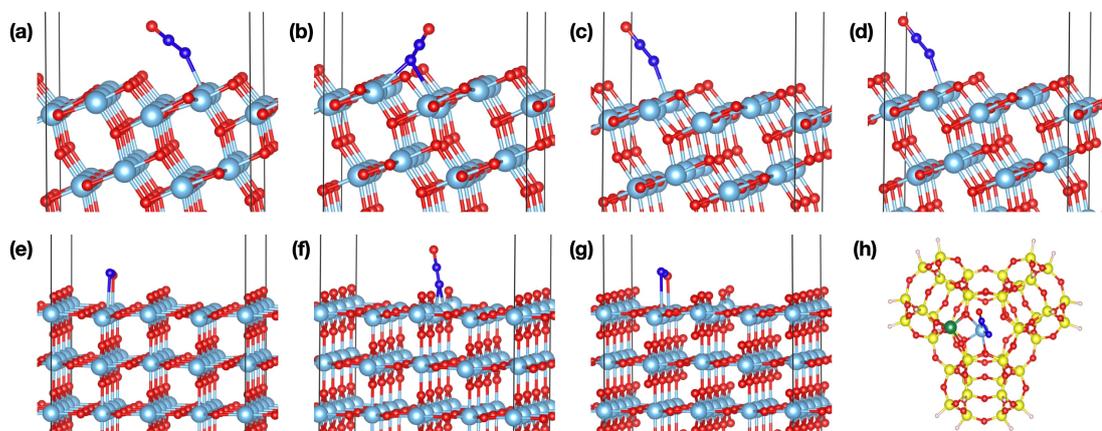
**Fig. C.1** The optimized structures of N adsorbed on (a)  $Ti_{5c}$  of perfect anatase  $TiO_2(101)$  surface, (b) oxygen vacancy( $Vo$ ) site of defective anatase  $TiO_2(101)$  surface, (c)  $Ti_{4c}$  site of defective anatase  $TiO_2(101)$  surface, (d)  $Ti_{5c}$  site of defective anatase  $TiO_2(101)$  surface, (e)  $Ti_{5c}$  of perfect rutile  $TiO_2(110)$  surface, (f) oxygen vacancy( $Vo$ ) site of defective rutile  $TiO_2(110)$  surface, (g)  $Ti_{4c}$  site of defective rutile  $TiO_2(110)$  surface, (d)  $Ti_{5c}$  site of defective rutile  $TiO_2(110)$  surface



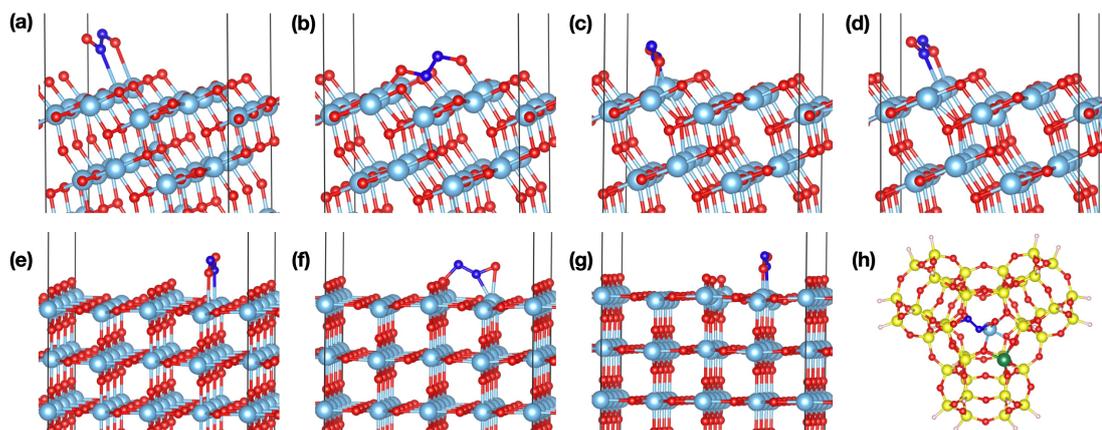
**Fig. C.2** The optimized structures of O adsorbed on (a)  $Ti_{5c}$  of perfect anatase  $TiO_2(101)$  surface, (b) oxygen vacancy( $Vo$ ) site of defective anatase  $TiO_2(101)$  surface, (c)  $Ti_{4c}$  site of defective anatase  $TiO_2(101)$  surface, (d)  $Ti_{5c}$  site of defective anatase  $TiO_2(101)$  surface, (e)  $Ti_{5c}$  of perfect rutile  $TiO_2(110)$  surface, (f) oxygen vacancy( $Vo$ ) site of defective rutile  $TiO_2(110)$  surface, (g)  $Ti_{5c}$  site of defective rutile  $TiO_2(110)$  surface, (h) Ti site of Ti-oxide Y-zeolite



**Fig. C.3** The optimized structures of  $O_2$  adsorbed on (a)  $Ti_{5c}$  of perfect anatase  $TiO_2(101)$  surface, (b) oxygen vacancy( $Vo$ ) site of defective anatase  $TiO_2(101)$  surface, (c)  $Ti_{4c}$  site of defective anatase  $TiO_2(101)$  surface, (d)  $Ti_{5c}$  site of defective anatase  $TiO_2(101)$  surface, (e)  $Ti_{5c}$  of perfect rutile  $TiO_2(110)$  surface, (f) oxygen vacancy( $Vo$ ) site of defective rutile  $TiO_2(110)$  surface, (g)  $Ti_{5c}$  site of defective rutile  $TiO_2(110)$  surface, (h) Ti site of Ti-oxide Y-zeolite



**Fig. C.4** The optimized structures of  $N_2O$  adsorbed on (a)  $Ti_{5c}$  of perfect anatase  $TiO_2(101)$  surface, (b) oxygen vacancy( $Vo$ ) site of defective anatase  $TiO_2(101)$  surface, (c)  $Ti_{4c}$  site of defective anatase  $TiO_2(101)$  surface, (d)  $Ti_{5c}$  site of defective anatase  $TiO_2(101)$  surface, (e)  $Ti_{5c}$  of perfect rutile  $TiO_2(110)$  surface, (f) oxygen vacancy( $Vo$ ) site of defective rutile  $TiO_2(110)$  surface, (g)  $Ti_{5c}$  site of defective rutile  $TiO_2(110)$  surface, (h)  $Ti$  site of  $Ti$ -oxide Y-zeolite



**Fig. C.5** The optimized structures of  $N_2O_2$  adsorbed on (a)  $Ti_{5c}$  of perfect anatase  $TiO_2(101)$  surface, (b) oxygen vacancy( $Vo$ ) site of defective anatase  $TiO_2(101)$  surface, (c)  $Ti_{4c}$  site of defective anatase  $TiO_2(101)$  surface, (d)  $Ti_{5c}$  site of defective anatase  $TiO_2(101)$  surface, (e)  $Ti_{5c}$  of perfect rutile  $TiO_2(110)$  surface, (f) oxygen vacancy( $Vo$ ) site of defective rutile  $TiO_2(110)$  surface, (g)  $Ti_{5c}$  site of defective rutile  $TiO_2(110)$  surface, (h)  $Ti$  site of  $Ti$ -oxide Y-zeolite